## ZEROS OF SCHROEDINGER EIGENFUNCTIONS AT POTENTIAL SINGULARITIES

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Let  $\Omega$  be a domain in  $\mathbb{R}^N$   $(N \geq 3)$ ,  $0 \in \Omega$  and let  $q \in L^1_{loc}(\Omega)$  denote a positive valued function. Recently Alimov and Joó [1], [2] achieved interesting results concerning the summability of eigenfunction expansions of the Schroedinger operator L defined as a self-adjoint extension of  $L_0 \colon C_0^{\infty}(\Omega) \ni f \mapsto -\Delta f + qf$  under the hypothesis  $q = q_0 + q_1$  where  $q_0 \in L^2(\Omega)$  is a radially symmetric positive function,  $q(x) \leq \text{const} \cdot |x|^{-2}$  and  $q_1 \in L^{\infty}(\Omega)$ . Their considerations are based upon Titchmarsh's mean value formula whose proof requires, as far we know (cf. [3]), the Gauss—Green theorem and hence the smoothness of the eigenfunctions of L. Nevertheless, the above conditions on q seem not too restrictive from the view points of applications e.g. in quantum chemistry, except for the assumption on the singularity of  $q_0$ . The aim of this short note is to point out that a bit stronger singularity of the potential q at 0 entails the vanishing of all the Schroedinger eigenfunctions at 0 even in a much more general setting. This fact may have interest in itself from the view points of quantum chemistry [4].

PROPOSITION. Suppose  $u \in C(\Omega)$  is a function satisfying the equation  $-\Delta u + qu = \lambda u$  in distribution sense for some  $\lambda$ . Then u(0) = 0 unless  $\overline{\lim}_{r \to 0} r^{-N+2} \int_{|x| < r} q(x) dx < \infty$ .

PROOF. Let  $u(0)\neq 0$ . Then we may assume without loss of generality u(0)=1 and we may choose  $r_0>0$  such that  $\{x\colon |x|< r_0\}\subset \Omega$  and  $\sup_{|x|< r_0}|u(x)-1|\leq 1/2$ . Let us fix  $r\in (0,r_0)$  arbitrarily and consider the function  $\psi_0(r)\equiv \mathrm{sgn}\left(\tau-\frac{3}{4}\,r\right)\chi(\tau)$  where  $\chi$  denotes the characteristic function of the interval  $\left(\frac{r}{2},\,r\right)$ . Observe that for a suitable sequence  $\psi_n\in C_0^\infty\left(\frac{r}{2},\,r\right),\ |\psi_n|\leq 1\ (n=1,2,\ldots)$  we have  $\psi_n(\tau)\to\psi_0(\tau)$   $(n\to\infty)$  whenever  $\tau\in [0,\infty)$ . Now define the functions  $\varphi_n\colon [0,\infty)\to \mathbb{R}$  and  $f_n\colon \Omega\to \mathbb{R}$  by

$$\varphi_n(\varrho) \equiv \int\limits_0^\infty \int\limits_{\xi}^\infty \psi_n(\tau) d\tau \, d\xi \quad \text{and} \quad f_n(x) \equiv \varphi_n(|x|) \quad (n = 0, 1, 2, \ldots).$$

Clearly,  $f_n \in C_0^{\infty}(\Omega)$  (n=1, 2, ...) and

$$\Delta f_n(x) = \varphi_n''(|x|) + \frac{N-1}{|x|} \varphi_n'(|x|) \to \Delta f_0(x) \quad (n \to \infty)$$

while

$$|\Delta f_n(x)| \le 1 + \frac{N-1}{|x|} \left| |x| - \frac{3}{4}r \right| \le \frac{N+1}{2} \quad \text{if} \quad \frac{r}{2} \le |x| \le r$$

and  $\Delta f_n(x) = 0$  elsewhere. Thus

$$\int_{\Omega} u[-\Delta f_n - \lambda f_n] \to \int_{\Omega} u[-\Delta f_0 - \lambda f_0]$$

and

$$\int_{O} qu f_{n} \to \int_{O} qu f_{0} \quad (n \to \infty).$$

By hypothesis,

$$\int_{\Omega} u(-\Delta f + qf) = \lambda \int_{\Omega} uf$$

for all  $f \in C_0^{\infty}(\Omega)$ , in particular for  $f = f_1, f_2, \ldots$  By passing to the limit, hence we obtain

(1) 
$$\int_{\Omega} u(\Delta f_0 + \lambda f_0) = \int_{\Omega} q u f_0.$$

However,

$$\left| \int_{\Omega} u(\Delta f_0 + \lambda f_0) \right| \leq \left( \frac{N+1}{2} + |\lambda| \right) \int_{r/2 < |x| < r} |u(x)| \, dx \leq \left( \frac{N+1}{2} + |\lambda| \right) \frac{3}{2} \, \omega_N r^N$$

where  $\omega_N$  denotes the volume of the unit ball in  $\mathbb{R}^N$ . On the other hand,

$$\operatorname{Re} \int_{\Omega} q u f_0 = \int_{\Omega} q(\operatorname{Re} u) f_0 \ge \int_{|x| < r/2} q(x) \frac{1}{2} \frac{r^2}{16} dx$$

since  $q, f_0 \ge 0$  and  $f_0(x) = r^2/16$  whenever |x| < r/2. Therefore, by (1),  $r^2 \cdot \int_{|x| < r/2} q \le cr^N$  where the constant c does not depend on  $r(< r_0)$ . This completes the proof.

COROLLARY. If  $q(x) \ge c|x|^{-(2+\epsilon)}$   $(x \in \Omega)$  for some  $c, \epsilon > 0$  then every continuous solution u of the Schroedinger equation  $-\Delta u + qu = \lambda u$  vanishes at 0.

## References

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(Received March 2, 1983)

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Acta Mathematica Hungarica 44, 1984