# Asymptotic behavior for minimizers of a p-energy functional associated with p-harmonic maps 

Yutian Lei<br>Depart. of Math., Nanjing Normal University, Nanjing 210097, P.R.China<br>E-mail: leiyutian@njnu.edu.cn


#### Abstract

The author studies the asymptotic behavior of minimizers $u_{\varepsilon}$ of a p-energy functional with penalization as $\varepsilon \rightarrow 0$. Several kinds of convergence for the minimizer to the p-harmonic map are presented under different assumptions.


Keywords: p-energy functional, p-energy minimizer, p-harmonic map
MSC 35B25, 35J70, 49K20, 58G18

## 1 Introduction

Let $G \subset R^{2}$ be a bounded and simply connected domain with smooth boundary $\partial G$, and $B_{1}=\left\{x \in R^{2}\right.$ or the complex plane $\left.C ; x_{1}^{2}+x_{2}^{2}<1\right\}$. Denote $S^{1}=$ $\left\{x \in R^{3} ; x_{1}^{2}+x_{2}^{2}=1, x_{3}=0\right\}$ and $S^{2}=\left\{x \in R^{3} ; x_{1}^{2}+x_{2}^{2}+x_{3}^{2}=1\right\}$. The vector value function can be denoted as $u=\left(u_{1}, u_{2}, u_{3}\right)=\left(u^{\prime}, u_{3}\right)$. Let $g=\left(g^{\prime}, 0\right)$ be a smooth map from $\partial G$ into $S^{1}$. Recall that the energy functional

$$
E_{\varepsilon}(u)=\frac{1}{2} \int_{G}|\nabla u|^{2} d x+\frac{1}{2 \varepsilon^{2}} \int_{G} u_{3}^{2} d x
$$

with a small parameter $\varepsilon>0$ was introduced in the study of some simplified model of high-energy physics, which controls the statics of planner ferromagnets and antiferromagnets (see [9] and [12]). The asymptotic behavior of minimizers of $E_{\varepsilon}(u)$ had been studied by Fengbo Hang and Fanghua Lin in [7]. When the term $\frac{u_{3}^{2}}{2 \varepsilon^{2}}$ replaced by $\frac{\left(1-|u|^{2}\right)^{2}}{4 \varepsilon^{2}}$ and $S^{2}$ replaced by $R^{2}$, the problem becomes the simplified model of the Ginzburg-Landau theory for superconductors and was well studied in many papers such as [1][2] and [13]. These works show that the properties of harmonic map with $S^{1}$-value can be studied via researching the minimizers of the functional with some penalization terms. Indeed, Y.Chen and M.Struwe used the penalty method to establish the global existence of partial regular weak solutions of the harmonic map flow (see [4] and [6]). M.Misawa studied the p-harmonic maps by using the same idea of the penalty method in
[11]. Now, the functional

$$
E_{\varepsilon}(u, G)=\frac{1}{p} \int_{G}|\nabla u|^{p} d x+\frac{1}{2 \varepsilon^{p}} \int_{G} u_{3}^{2} d x, \quad p>2
$$

which equipped with the penalization $\frac{1}{2 \varepsilon^{p}} \int_{G} u_{3}^{2} d x$, will be considered in this paper. From the direct method in the calculus of variations, it is easy to see that the functional achieves its minimum in the function class $W_{g}^{1, p}\left(G, S^{2}\right)$. Without loss of generality, we assume $u_{3} \geq 0$, otherwise we may consider $\left|u_{3}\right|$ in view of the expression of the functional. We will research the asymptotic properties of minimizers of this p-energy functional on $W_{g}^{1, p}\left(G, S^{2}\right)$ as $\varepsilon \rightarrow 0$, and shall prove the limit of the minimizers is the p-harmonic map.

Theorem 1.1 Let $u_{\varepsilon}$ be a minimizer of $E_{\varepsilon}(u, G)$ on $W_{g}^{1, p}\left(G, S^{2}\right)$. Assume $\operatorname{deg}\left(g^{\prime}, \partial G\right)=0$. Then

$$
\lim _{\varepsilon \rightarrow 0} u_{\varepsilon}=\left(u_{p}, 0\right), \quad \text { in } \quad W^{1, p}\left(G, S^{2}\right)
$$

where $u_{p}$ is the minimizer of $\int_{G}|\nabla u|^{p} d x$ in $W_{g}^{1, p}\left(G, \partial B_{1}\right)$.

Remark. When $p=2$, [7] shows that if $\operatorname{deg}\left(g^{\prime}, \partial G\right)=0$, the minimizer of $E_{\varepsilon}(u)$ in $H_{g}^{1}\left(G, S^{2}\right)$ is just $\left(u_{2}, 0\right)$, where $u_{2}$ is the energy minimizer, i.e., it is the minimizer of $\int_{G}|\nabla u|^{2} d x$ in $H_{g}^{1}\left(G, \partial B_{1}\right)$. When $p>2$, there may be several minimizers of $E_{\varepsilon}(u, G)$ in $W_{g}^{1, p}\left(G, S^{2}\right)$. The author proved that there exists a minimizer, which is called the regularized minimizer, is just $\left(u_{p}, 0\right)$, where $u_{p}$ is the minimizer of $\int_{G}|\nabla u|^{p} d x$ in $W_{g}^{1, p}\left(G, \partial B_{1}\right)$. For the other minimizers, we only deduced the result as Theorem 1.1.

Comparing with the assumption of Theorem 1.1, we will consider the problem under some weaker conditions. Then we have

Theorem 1.2 Assume $u_{\varepsilon}$ is a critical point of $E_{\varepsilon}(u, G)$ on $W_{g}^{1, p}\left(G, S^{2}\right)$. If

$$
\begin{equation*}
E_{\varepsilon}\left(u_{\varepsilon}, K\right) \leq C \tag{1.1}
\end{equation*}
$$

for some subdomain $K \subseteq G$. Then there exists a subsequence $u_{\varepsilon_{k}}$ of $u_{\varepsilon}$ such that as $k \rightarrow \infty$,

$$
\begin{equation*}
u_{\varepsilon_{k}} \rightarrow\left(u_{p}, 0\right), \quad \text { weakly in } \quad W^{1, p}\left(K, R^{3}\right), \tag{1.2}
\end{equation*}
$$

where $u_{p}$ is a critical point of $\int_{K}|\nabla u|^{p} d x$ in $W^{1, p}\left(K, \partial B_{1}\right)$, which is named p-harmonic map on $K$. Moreover, for any $\zeta \in C_{0}^{\infty}(K)$, when $\varepsilon \rightarrow 0$,

$$
\begin{equation*}
\int_{K}\left|\nabla u_{\varepsilon_{k}}\right|^{p} \zeta d x \rightarrow \int_{K}\left|\nabla u_{p}\right|^{p} \zeta d x \tag{1.3}
\end{equation*}
$$

$$
\begin{equation*}
\frac{1}{\varepsilon_{k}^{p}} \int_{K} u_{\varepsilon_{k} 3} \zeta d x \rightarrow 0 \tag{1.4}
\end{equation*}
$$

The convergent rate of $\left|u_{\varepsilon}^{\prime}\right| \rightarrow 1$ and $u_{3} \rightarrow 0$ will be concerned with as $\varepsilon \rightarrow 0$.
Theorem 1.3 Let $u_{\varepsilon}$ be a minimizer of $E_{\varepsilon}(u, G)$ on $W_{g}^{1, p}\left(G, S^{2}\right)$. If (1.1) holds, then there exists a positive constant $C$, such that as $\varepsilon \rightarrow 0$,

$$
\left.\int_{K}|\nabla| u_{\varepsilon}^{\prime}\right|^{p} d x+\int_{K}\left|\nabla u_{\varepsilon 3}\right|^{p} d x+\frac{1}{\varepsilon^{p}} \int_{K} u_{\varepsilon 3}^{2} d x \leq C \varepsilon^{\beta},
$$

where $\beta=1-\frac{2}{p}$ when $p \in\left(2, p_{0}\right] ; \beta=\frac{2 p}{p^{2}-2}$ when $p>p_{0}$. Here $p_{0} \in(4,5)$ is a constant satisfying $p^{3}-4 p^{2}-2 p+4=0$.

## 2 Proof of Theorem 1.1

In this section, we always assume $\operatorname{deg}\left(g^{\prime}, \partial G\right)=0$. By the argument of the weak low semi-continuity, it is easy to deduce the strong convergence in $W^{1, p}$ sense for some subsequence of the minimizer $u_{\varepsilon}$. To improve the conclusion of the convergence for all $u_{\varepsilon}$, we need to research the limit function: p-harmonic map.

From $\operatorname{deg}\left(g^{\prime}, \partial G\right)=0$ and the smoothness of $\partial G$ and $g$, we see that there is a smooth function $\phi_{0}: \partial G \rightarrow R$ such that

$$
\begin{equation*}
g=e^{i \phi_{0}}, \quad \text { on } \quad \partial G . \tag{2.1}
\end{equation*}
$$

Consider the Dirichlet problem

$$
\begin{gather*}
-\operatorname{div}\left(|\nabla \Phi|^{p-2} \nabla \Phi\right)=0, \quad \text { in } \quad G  \tag{2.2}\\
\left.\Phi\right|_{\partial G}=\phi_{0} \tag{2.3}
\end{gather*}
$$

Proposition 2.1 There exists the unique weak solution $\Phi$ of (2.2) and (2.3) in $W^{1, p}(G, R)$. Namely, for any $\phi \in W_{0}^{1, p}(G, R)$, there is the unique $\Phi$ satisfies

$$
\begin{equation*}
\int_{G}|\nabla \Phi|^{p-2} \nabla \Phi \nabla \phi d x=0 \tag{2.4}
\end{equation*}
$$

Proof. By using the method in the calculus of variations, we can see the existence for the weak solution of (2.2) and (2.3).

If both $\Phi_{1}$ and $\Phi_{2}$ are weak solutions of (2.2) and (2.3), then, by taking the test function $\phi=\Phi_{1}-\Phi_{2}$ in (2.4), there holds

$$
\int_{G}\left(\left|\nabla \Phi_{1}\right|^{p-2} \nabla \Phi_{1}-\left|\nabla \Phi_{2}\right|^{p-2} \nabla \Phi_{2}\right) \nabla\left(\Phi_{1}-\Phi_{2}\right) d x=0
$$

In view of Lemma 1.2 in [5] we have

$$
\int_{G}\left|\nabla\left(\Phi_{1}-\Phi_{2}\right)\right|^{p} d x \leq 0 .
$$

Hence, $\Phi_{1}-\Phi_{2}=$ Const. on $\bar{G}$. Noting the boundary condition, we see $\Phi_{1}-\Phi_{2}=$ 0 on $\bar{G}$. Proposition is proved.

Recall that $u \in W_{g}^{1, p}\left(G, \partial B_{1}\right)$ is named p-harmonic map, if it is the critical point of $\int_{G}|\nabla u|^{p} d x$. Namely, it is the weak solution of

$$
\begin{equation*}
-\operatorname{div}\left(|\nabla u|^{p-2} \nabla u\right)=u|\nabla u|^{p} \tag{2.5}
\end{equation*}
$$

on $G$, or for any $\phi \in C_{0}^{\infty}\left(G, R^{2}\right.$ or $\left.C\right)$, it satisfies

$$
\begin{equation*}
\int_{G}|\nabla u|^{p-2} \nabla u \nabla \phi d x=\int_{G} u|\nabla u|^{p} \phi d x . \tag{2.6}
\end{equation*}
$$

Assume $\Phi$ is the unique weak solution of (2.2) and (2.3). Set

$$
\begin{equation*}
u_{p}=e^{i \Phi}, \quad \text { on } \quad \bar{G} . \tag{2.7}
\end{equation*}
$$

Proposition $2.2 u_{p}$ defined in (2.7) is a $p$-harmonic map on $G$.
Proof. Obviously, $u_{p} \in W_{g}^{1, p}\left(G, \partial B_{1}\right)$ since $\Phi \in W_{\phi_{0}}^{1, p}(G, R)$. We only need to prove that $u_{p}$ satisfies (2.6) for any $\phi \in C_{0}^{\infty}(G, C)$. In fact,

$$
\begin{aligned}
& \int_{G}\left(\left|\nabla u_{p}\right|^{p-2} \nabla u_{p} \nabla \phi-u_{p} \phi\left|\nabla u_{p}\right|^{p}\right) d x \\
= & i \int_{G}|\nabla \Phi|^{p-2} \nabla \Phi\left(e^{i \Phi} \nabla \phi+i e^{i \Phi} \nabla \Phi \phi\right) d x=i \int_{G}|\nabla \Phi|^{p-2} \nabla \Phi \nabla\left(e^{i \Phi} \phi\right) d x
\end{aligned}
$$

for any $\phi \in C_{0}^{\infty}(G, C)$. Noting $e^{i \Phi} \phi \in W_{0}^{1, p}(G, C)$ and $\Phi$ is the weak solution of (2.2) and (2.3), we obtain

$$
\int_{G}\left|\nabla u_{p}\right|^{p-2} \nabla u_{p} \nabla \phi d x-\int_{G} u_{p} \phi\left|\nabla u_{p}\right|^{p} d x=0
$$

for any $\phi \in C_{0}^{\infty}(G, C)$. Proposition is proved.
Since $W_{g}^{1, p}\left(G, \partial B_{1}\right) \neq \emptyset$ when $\operatorname{deg}\left(g^{\prime}, \partial G\right)=0$, we may consider the minimization problem

$$
\begin{equation*}
\operatorname{Min}\left\{\int_{G}|\nabla u|^{p} d x ; u \in W_{g}^{1, p}\left(G, \partial B_{1}\right)\right\} \tag{2.8}
\end{equation*}
$$

The solution is called p-energy minimizer.
Proposition 2.3 The solution of (2.8) exists.

Proof. The weakly low semi-continuity of $\int_{G}|\nabla u|^{p} d x$ is well-known. On the other hand, if taking a minimizing sequence $u_{k}$ of $\int_{G}|\nabla u|^{p} d x$ in $W_{g}^{1, p}\left(G, \partial B_{1}\right)$, then there is a subsequence of $u_{k}$, which is still denoted $u_{k}$ itself, such that as $k \rightarrow \infty, u_{k}$ converges to $u_{0}$ weakly in $W^{1, p}(G, C)$. Noting that $W_{g}^{1, p}\left(G, \partial B_{1}\right)$ is the weakly closed subset of $W^{1, p}(G, C)$ since it is the convex closed subset, we see that $u_{0} \in W_{g}^{1, p}\left(G, \partial B_{1}\right)$. Thus, if denote

$$
\alpha=\operatorname{Inf}\left\{\int_{G}|\nabla u|^{p} d x ; u \in W_{g}^{1, p}\left(G, \partial B_{1}\right)\right\},
$$

then

$$
\alpha \leq \int_{G}\left|\nabla u_{0}\right|^{p} d x \leq \underline{\lim }_{k \rightarrow \infty} \int_{G}\left|\nabla u_{k}\right|^{p} d x \leq \alpha .
$$

This means $u_{0}$ is the solution of (2.8).
Obviously, the p-energy minimizer is the p-harmonic map.
Proposition 2.4 The p-harmonic map is unique in $W_{g}^{1, p}\left(G, \partial B_{1}\right)$.
Proof. It follows that $u_{p}=e^{i \Phi}$ is a p-harmonic map from Proposition 2.2. If $u$ is also a p-harmonic map in $W_{g}^{1, p}\left(G, \partial B_{1}\right)$, then from $\operatorname{deg}\left(g^{\prime}, \partial G\right)=0$ and using the results in [3], we know that there is $\Phi_{0} \in W^{1, p}(G, R)$ such that

$$
\begin{gathered}
u=e^{i \Phi_{0}}, \quad \text { on } \quad \bar{G}, \\
\Phi_{0}=\phi_{0}, \quad \text { on } \quad \partial G .
\end{gathered}
$$

Substituting these into (2.6), we see that $\Phi_{0}$ is a weak solution of (2.2) and (2.3). Proposition 2.1 leads to $\Phi_{0}=\Phi$, which implies $u=u_{p}$.

Now, we conclude that $u_{0}$ in Proposition 2.3 is just the p-harmonic map $u_{p}$. Furthermore, the p-energy minimizer is also unique in $W_{g}^{1, p}\left(G, \partial B_{1}\right)$.

Proof of Theorem 1.1. Noticing that $u_{\varepsilon}$ is the minimizer, we have

$$
\begin{equation*}
E_{\varepsilon}\left(u_{\varepsilon}, G\right) \leq E_{\varepsilon}\left(\left(u_{p}, 0\right), G\right) \leq C \tag{2.9}
\end{equation*}
$$

with $C>0$ independent of $\varepsilon$. This means

$$
\begin{align*}
& \int_{G}\left|\nabla u_{\varepsilon}\right|^{p} d x \leq C  \tag{2.10}\\
& \int_{G} u_{\varepsilon 3}^{2} d x \leq C \varepsilon^{p} \tag{2.11}
\end{align*}
$$

Using (2.10), $\left|u_{\varepsilon}\right|=1$ and the embedding theorem, we see that there exists a subsequence $u_{\varepsilon_{k}}$ of $u_{\varepsilon}$ and $u_{*} \in W^{1, p}\left(G, R^{3}\right)$, such that as $\varepsilon_{k} \rightarrow 0$,

$$
\begin{gather*}
u_{\varepsilon_{k}} \rightarrow u_{*}, \quad \text { weakly in } \quad W^{1, p}\left(G, S^{2}\right),  \tag{2.12}\\
u_{\varepsilon_{k}} \rightarrow u_{*}, \quad \text { in } \quad C^{\alpha}\left(\bar{G}, S^{2}\right), \quad \alpha \in(0,1-2 / p) . \tag{2.13}
\end{gather*}
$$

Obviously, (2.11) and (2.13) lead to $u_{*} \in W_{g}^{1, p}\left(G, S^{1}\right)$.
Applying (2.12) and the weak low semi-continuity of $\int_{G}|\nabla u|^{p} d x$, we have

$$
\int_{G}\left|\nabla u_{*}\right|^{p} d x \leq \underline{\lim }_{\varepsilon_{k} \rightarrow 0} \int_{G}\left|\nabla u_{\varepsilon_{k}}\right|^{p} d x
$$

On the other hand, (2.9) implies

$$
\int_{G}\left|\nabla u_{\varepsilon_{k}}\right|^{p} d x \leq \int_{G}\left|\nabla\left(u_{p}, 0\right)\right|^{p} d x
$$

hence,

$$
\int_{G}\left|\nabla u_{*}^{\prime}\right|^{p} d x \leq \int_{G}\left|\nabla u_{p}\right|^{p} d x .
$$

This means that $u_{*}^{\prime}$ is also a p-energy minimizer. Noting the uniqueness we see $u_{*}=u_{p}$. Thus

$$
\int_{G}\left|\nabla u_{p}\right|^{p} d x \leq \underline{\lim }_{\varepsilon_{k} \rightarrow 0} \int_{G}\left|\nabla u_{\varepsilon_{k}}\right|^{p} d x \leq \overline{\lim }_{\varepsilon_{k} \rightarrow 0} \int_{G}\left|\nabla u_{\varepsilon_{k}}\right|^{p} d x \leq \int_{G}\left|\nabla u_{p}\right|^{p} d x .
$$

When $\varepsilon_{k} \rightarrow 0$,

$$
\int_{G}\left|\nabla u_{\varepsilon_{k}}\right|^{p} \rightarrow \int_{G}\left|\nabla u_{p}\right|^{p} .
$$

Combining this with (2.12) yields

$$
\lim _{k \rightarrow \infty} \nabla u_{\varepsilon_{k}}=\nabla\left(u_{p}, 0\right), \quad \text { in } \quad L^{p}\left(G, S^{2}\right)
$$

In addition, (2.13) implies that as $\varepsilon \rightarrow 0$,

$$
u_{\varepsilon_{k}} \rightarrow\left(u_{p}, 0\right), \quad \text { in } \quad L^{p}\left(G, S^{2}\right)
$$

Then

$$
\lim _{k \rightarrow \infty} u_{\varepsilon_{k}}=\left(u_{p}, 0\right), \quad \text { in } \quad W^{1, p}\left(G, S^{2}\right) .
$$

Noticing the uniqueness of $\left(u_{p}, 0\right)$, we see the convergence above also holds for all $u_{\varepsilon}$.

## 3 Proof of Theorem 1.2

In this section, we always assume that $u_{\varepsilon}$ is the critical point of the functional, and $E_{\varepsilon}\left(u_{\varepsilon}, K\right) \leq C$ for some subdomain $K \subseteq G$, where $C$ is independent of $\varepsilon$. The assumption is weaker than that of Theorem 1.1. So, all the results in this section will be derived in the weak sense.

The method in the calculus of variations shows that the minimizer $u_{\varepsilon} \in$ $W_{g}^{1, p}\left(G, S^{2}\right)$ is a weak solution of

$$
\begin{equation*}
-\operatorname{div}\left(|\nabla u|^{p-2} \nabla u\right)=u|\nabla u|^{p}+\frac{1}{\varepsilon^{p}}\left(u u_{3}^{2}-u_{3} e_{3}\right), \quad \text { on } \quad G, \tag{3.1}
\end{equation*}
$$

where $e_{3}=(0,0,1)$. Namely, for any $\psi \in W_{0}^{1, p}\left(G, R^{3}\right), u_{\varepsilon}$ satisfies

$$
\begin{equation*}
\int_{G}|\nabla u|^{p-2} \nabla u \nabla \psi d x=\int_{G} u \psi|\nabla u|^{p} d x+\frac{1}{\varepsilon^{p}} \int_{G} \psi\left(u u_{3}^{2}-u_{3} e_{3}\right) d x . \tag{3.2}
\end{equation*}
$$

Proof of (1.2). $E_{\varepsilon}\left(u_{\varepsilon}, K\right) \leq C$ means

$$
\begin{align*}
& \int_{K}\left|\nabla u_{\varepsilon}\right|^{p} d x \leq C,  \tag{3.3}\\
& \int_{K} u_{\varepsilon 3}^{2} d x \leq C \varepsilon^{p} \tag{3.4}
\end{align*}
$$

where $C$ is independent of $\varepsilon$. Combining the fact $\left|u_{\varepsilon}\right|=1$ a.e. on $\bar{G}$ with (3.3) we know that there exist $u_{p} \in W^{1, p}\left(K, \partial B_{1}\right)$ and a subsequence $u_{\varepsilon_{k}}$ of $u_{\varepsilon}$, such that as $\varepsilon_{k} \rightarrow 0$,

$$
\begin{gather*}
u_{\varepsilon_{k}} \rightarrow\left(u_{p}, 0\right), \quad \text { weakly in } \quad W^{1, p}(K),  \tag{3.5}\\
u_{\varepsilon_{k}} \rightarrow\left(u_{p}, 0\right), \quad \text { in } \quad C^{\alpha}(\bar{K}), \tag{3.6}
\end{gather*}
$$

for some $\alpha \in\left(0,1-\frac{2}{p}\right)$. In the following we will prove that $u_{p}$ is a weak solution of (2.5).

Let $B=B(x, 3 R) \subset \subset K . \phi \in C_{0}^{\infty}(B(x, 3 R) ;[0,1]), \phi=1$ on $B(x, R), \phi=0$ on $B \backslash B(x, 2 R)$ and $|\nabla \phi| \leq C$, where $C$ is independent of $\varepsilon$. Denote $u=u_{\varepsilon_{k}}$ in (3.2) and take $\psi=(0,0, \phi)$. Thus

$$
\int_{B}|\nabla u|^{p-2} \nabla u_{3} \nabla \phi d x+\frac{1}{\varepsilon_{k}^{p}} \int_{B}\left|u^{\prime}\right|^{2} \phi u_{3} d x=\int_{B} u_{3} \phi|\nabla u|^{p} d x .
$$

Applying (3.3) we can derive that

$$
\begin{equation*}
\frac{1}{\varepsilon_{k}^{p}} \int_{B}\left|u^{\prime}\right|^{2} \phi\left|u_{3}\right| d x \leq \int_{B}|\nabla u|^{p} \phi d x+\int_{B}|\nabla u|^{p-1}|\nabla \phi| d x \leq C . \tag{3.7}
\end{equation*}
$$

From (3.6) it follows $\left|u^{\prime}\right| \geq 1 / 2$ when $\varepsilon_{k}$ is sufficiently small. Noting $\phi=1$ on $B(x, R)$, we have

$$
\begin{equation*}
\frac{1}{\varepsilon_{k}^{p}} \int_{B(x, R)}\left|u_{3}\right| d x \leq C . \tag{3.8}
\end{equation*}
$$

Taking $\frac{1}{k}=\varepsilon_{k}, F_{k}=\frac{1}{\varepsilon_{k}^{p}}\left(u_{\varepsilon_{k}} u_{\varepsilon_{k} 3}^{2}-u_{\varepsilon_{k} 3} e_{3}\right)$ in Lemma 3.11 of [8], noting $\left|F_{k}\right|=$ $\frac{1}{\varepsilon_{k}^{p}}\left|u_{3}\right|\left|u^{\prime}\right|$ and applying (3.5) and (3.8) we obtain that for any $q \in(1, p)$, as $\varepsilon_{k} \rightarrow 0, \nabla u_{\varepsilon_{k}} \rightarrow \nabla u_{p}$, in $L^{q}(B(x, R))$. Since $B(x, R)$ is an arbitrary disc in $K$, we can see that as $\varepsilon_{k} \rightarrow 0$, for any $\xi \in C_{0}^{\infty}\left(B, R^{3}\right)$ there holds

$$
\begin{equation*}
\int_{B}\left|\nabla u_{\varepsilon_{k}}\right|^{p-2} \nabla u_{\varepsilon_{k}} \nabla \xi d x \rightarrow \int_{B}\left|\nabla u_{p}\right|^{p-2} \nabla u_{p} \nabla \xi d x . \tag{3.9}
\end{equation*}
$$

Now, denote $u^{\prime}=u_{\varepsilon_{k}}^{\prime}=\left(u_{1}, u_{2}\right)$. Taking $\psi=\left(u_{2}, 0,0\right) \zeta$ and $\psi=\left(0, u_{1}, 0\right) \zeta$ in (3.2), respectively, where $\zeta \in C_{0}^{\infty}(B, R)$, we have that for $m, j \in\{1,2\}$, and $m \neq j$,

$$
\begin{aligned}
& \frac{1}{\varepsilon_{k}^{p}} \int_{B} u_{3}^{2} u_{m} u_{j} \zeta d x+\int_{B} u_{m} u_{j} \zeta|\nabla u|^{p} d x \\
& =\int_{B}|\nabla u|^{p-2} \nabla u_{m} \nabla u_{j} \zeta d x+\int_{B} u_{j}|\nabla u|^{p-2} \nabla u_{m} \nabla \zeta d x .
\end{aligned}
$$

One equation subtracts the other one, then

$$
\begin{equation*}
0=\int_{B}|\nabla u|^{p-2}(u \wedge \nabla u) \nabla \zeta d x \tag{3.10}
\end{equation*}
$$

where $u \wedge \nabla u=u_{1} \nabla u_{2}-u_{2} \nabla u_{1}$. On the other hand, since

$$
\begin{aligned}
& \int_{B} u_{2}|\nabla u|^{p-2} \nabla u_{1} \nabla \zeta d x-\int_{B} u_{p 2}\left|\nabla u_{p}\right|^{p-2} \nabla u_{p 1} \nabla \zeta d x \\
= & \int_{B}\left(|\nabla u|^{p-2} \nabla u_{1}-\left|\nabla u_{p}\right|^{p-2} \nabla u_{p 1}\right) u_{p 2} \nabla \zeta d x \\
& +\int_{B}|\nabla u|^{p-2} \nabla u_{1} \nabla \zeta\left(u_{2}-u_{p 2}\right) d x,
\end{aligned}
$$

we obtain that as $\varepsilon_{k} \rightarrow 0$,

$$
\begin{equation*}
\int_{B} u_{2}|\nabla u|^{p-2} \nabla u_{1} \nabla \zeta d x \rightarrow \int_{B} u_{p 2}\left|\nabla u_{p}\right|^{p-2} \nabla u_{p 1} \nabla \zeta d x \tag{3.11}
\end{equation*}
$$

by using (3.3)(3.6) and (3.9). Similarly, we may also get that

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \int_{B} u_{1}|\nabla u|^{p-2} \nabla u_{2} \nabla \zeta d x=\int_{B} u_{p 1}\left|\nabla u_{p}\right|^{p-2} \nabla u_{p 2} \nabla \zeta d x . \tag{3.12}
\end{equation*}
$$

(3.12) subtracts (3.11), then

$$
\lim _{\varepsilon \rightarrow 0} \int_{B}|\nabla u|^{p-2}(u \wedge \nabla u) \nabla \zeta d x=\int_{B}\left|\nabla u_{p}\right|^{p-2}\left(u_{p} \wedge \nabla u_{p}\right) \nabla \zeta d x .
$$

Combining this with (3.10), we have

$$
\begin{equation*}
\int_{B}\left|\nabla u_{p}\right|^{p-2}\left(u_{p} \wedge \nabla u_{p}\right) \nabla \zeta d x=0 \tag{3.13}
\end{equation*}
$$

Let $u_{*}=u_{p 1}+i u_{p 2}: B \rightarrow C$. Thus

$$
\begin{equation*}
\left|\nabla u_{*}\right|^{2}=\left|\nabla u_{p}\right|^{2} \tag{3.14}
\end{equation*}
$$

It is easy to see that $\overline{u_{*}} \nabla u_{*}=\nabla\left(\left|u_{*}\right|^{2}\right)+\left(u_{*} \wedge \nabla u_{*}\right) i=0+\left(u_{*} \wedge \nabla u_{*}\right) i$ since $\left|u_{*}\right|^{2}=\left|u_{p 1}\right|^{2}+\left|u_{p 2}\right|^{2}=1$. Substituting this into (3.13) yields

$$
-i \int_{B}\left|\nabla u_{*}\right|^{p-2} \overline{u_{*}} \nabla u_{*} \nabla \zeta d x=0
$$

for any $\zeta \in C_{0}^{\infty}(B, R)$. Taking $\zeta=\operatorname{Re}\left(u_{*} \phi_{j}\right)$ and $\zeta=\operatorname{Im}\left(u_{*} \phi_{j}\right)(j=1,2)$, respectively, where $\phi=\left(\phi_{1}, \phi_{2}\right) \in C_{0}^{\infty}\left(B, R^{2}\right)$, we can see that

$$
\int_{B}\left|\nabla u_{*}\right|^{p-2} \overline{u_{*}} \nabla u_{*} \nabla \operatorname{Re}\left(u_{*} \phi\right) d x+i \int_{B}\left|\nabla u_{*}\right|^{p-2} \overline{u_{*}} \nabla u_{*} \nabla \operatorname{Im}\left(u_{*} \phi\right) d x=0
$$

Namely

$$
0=\int_{G}\left|\nabla u_{*}\right|^{p-2} \overline{u_{*}} \nabla u_{*} \nabla\left(u_{*} \phi\right) d x .
$$

Noting $\overline{u_{*}} \nabla u_{*}=-u_{*} \nabla \overline{u_{*}}$, we obtain

$$
\begin{aligned}
0 & =\int_{B}\left|\nabla u_{*}\right|^{p-2} \nabla u_{*} \nabla \phi d x-\int_{B}\left|\nabla u_{*}\right|^{p-2} u_{*} \nabla \overline{u_{*}} \nabla u_{*} \phi d x \\
& =\int_{B}\left|\nabla u_{*}\right|^{p-2} \nabla u_{*} \nabla \phi d x-\int_{B}\left|\nabla u_{*}\right|^{p} u_{*} \phi d x:=J
\end{aligned}
$$

By using (3.14) and $\operatorname{Re}(J)=0, \operatorname{Im}(J)=0$, we have

$$
\begin{equation*}
\int_{B}\left|\nabla u_{p}\right|^{p-2} \nabla u_{p 1} \nabla \phi d x=\int_{B}\left|\nabla u_{p}\right|^{p} u_{p 1} \phi d x \tag{3.15}
\end{equation*}
$$

and

$$
\int_{B}\left|\nabla u_{p}\right|^{p-2} \nabla u_{p 2} \nabla \phi d x=\int_{B}\left|\nabla u_{p}\right|^{p} u_{p 2} \phi d x .
$$

Combining this with (3.15) yields that for any $\phi \in C_{0}^{\infty}\left(B, R^{3}\right)$,

$$
\int_{B}\left|\nabla u_{p}\right|^{p-2} \nabla u_{p} \nabla \phi d x=\int_{B}\left|\nabla u_{p}\right|^{p} u_{p} \phi d x
$$

It shows that $u_{p}$ is a weak solution of (2.5). (1.2) is completed.

Proof of (1.3). For simplification, denote $\varepsilon_{k}=\varepsilon$. From (3.3) and (3.6) it is deduced that as $\varepsilon \rightarrow 0$,

$$
\begin{align*}
& \left.\left.\quad\left|\int_{K} u_{3}^{2} \zeta\right| \nabla u\right|^{p} d x\left|\leq \sup _{K}\left(1-\left|u^{\prime}\right|^{2}\right) \cdot \int_{K}\right| \nabla u\right|^{p} d x \rightarrow 0,  \tag{3.16}\\
& \left.\left|\int_{K} u^{\prime} u_{p} \zeta\right| \nabla u\right|^{p} d x-\int_{K} \zeta|\nabla u|^{p} d x\left|=\left|\int_{K}\left(u^{\prime} u_{p}-u_{p} u_{p}\right) \zeta\right| \nabla u\right|^{p} d x \mid  \tag{3.17}\\
& \leq\left.\sup _{K}\left|u^{\prime}-u_{p}\right| \cdot\left|\int_{K} u_{p}\right| \nabla u\right|^{p} d x \mid \rightarrow 0,
\end{align*}
$$

and

$$
\begin{equation*}
\int_{K}\left(u-\left(u_{p}, 0\right)\right) \zeta|\nabla u|^{p} d x \leq\left.\sup _{K}\left|u-\left(u_{p}, 0\right)\right| \cdot\left|\int_{K} u_{p}\right| \nabla u\right|^{p} d x \mid \rightarrow 0 \tag{3.18}
\end{equation*}
$$

Similarly, (3.4) and (3.6) imply that as $\varepsilon \rightarrow 0$,

$$
\begin{equation*}
\left|\frac{1}{\varepsilon^{p}} \int_{K} u_{3}^{2} \zeta d x-\frac{1}{\varepsilon^{p}} \int_{K} u_{3}^{2} \zeta\left(1-u_{3}^{2}\right) d x\right| \leq \sup _{K}\left|1-\left|u^{\prime}\right|^{2}\right| \cdot \frac{1}{\varepsilon^{p}}\left|\int_{K} u_{3}^{2} d x\right| \rightarrow 0 \tag{3.19}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|\frac{1}{\varepsilon^{p}} \int_{K} u_{p} \zeta u^{\prime} u_{3}^{2} d x-\frac{1}{\varepsilon^{p}} \int_{K} \zeta u_{3}^{2} d x\right| \leq \sup _{K}\left|u^{\prime}-u_{p}\right| \cdot \frac{1}{\varepsilon^{p}}\left|\int_{K} u_{p} u_{3}^{2} d x\right| \rightarrow 0 . \tag{3.20}
\end{equation*}
$$

Letting $\varepsilon \rightarrow 0$ in (3.2) we have

$$
\begin{align*}
& \lim _{\varepsilon \rightarrow 0}\left[\int_{K} u \psi|\nabla u|^{p} d x+\frac{1}{\varepsilon^{p}} \int_{K} \psi\left(u u_{3}^{2}-u_{3} e_{3}\right) d x\right]  \tag{3.21}\\
= & \int_{K}\left|\nabla u_{p}\right|^{p-2} \nabla\left(u_{p}, 0\right) \nabla \psi d x=\int_{G}\left(u_{p}, 0\right) \psi\left|\nabla u_{p}\right|^{p} d x .
\end{align*}
$$

Take $\psi=\left(0,0, u_{3} \zeta\right)$ where $\zeta \in C_{0}^{\infty}(K)$ we have

$$
\lim _{\varepsilon \rightarrow 0}\left[\int_{K} u_{3}^{2} \zeta|\nabla u|^{p} d x+\frac{1}{\varepsilon^{p}} \int_{K} u_{3}^{2} \zeta\left(u_{3}^{2}-1\right) d x\right]=0
$$

Combining this with (3.16) we derive

$$
\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon^{p}} \int_{K} u_{3}^{2} \zeta\left(u_{3}^{2}-1\right) d x=0
$$

Substituting this into (3.19) yields

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon^{p}} \int_{K} u_{3}^{2} \zeta d x=0 \tag{3.22}
\end{equation*}
$$

Hence, as $\varepsilon \rightarrow 0$,

$$
\frac{1}{\varepsilon^{p}}\left|\int_{K} u u_{3}^{2} \zeta d x\right| \leq \frac{1}{\varepsilon^{p}} \int_{K} u_{3}^{2} \zeta d x \rightarrow 0
$$

Thus, for any $\psi \in W_{0}^{1, p}\left(K, R^{3}\right)$, there holds

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon^{p}} \int_{K} u u_{3}^{2} \psi d x=0 \tag{3.23}
\end{equation*}
$$

In addition, substituting (3.22) into (3.20) leads to

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon^{p}} \int_{K} u_{p} \zeta u^{\prime} u_{3}^{2} d x=0 . \tag{3.24}
\end{equation*}
$$

Take $\psi=\left(u_{p} \zeta, 0\right)$ in (3.21) we have

$$
\lim _{\varepsilon \rightarrow 0}\left[\int_{K} u^{\prime} u_{p} \zeta|\nabla u|^{p} d x+\frac{1}{\varepsilon^{p}} \int_{K} u_{p} \zeta u^{\prime} u_{3}^{2} d x\right]=\int_{K}\left|\nabla u_{p}\right|^{p} \zeta d x,
$$

which, together with (3.24), implies

$$
\lim _{\varepsilon \rightarrow 0} \int_{K} u^{\prime} u_{p} \zeta|\nabla u|^{p} d x=\int_{K}\left|\nabla u_{p}\right|^{p} \zeta d x
$$

Combining this with (3.17) we can see (1.3) at last.

Proof of (1.4). Obviously, (3.18) and (1.3) show that as $\varepsilon \rightarrow 0$,

$$
\begin{aligned}
& \left.\left|\int_{K} u\right| \nabla u\right|^{p} \psi d x-\int_{K}\left(u_{p}, 0\right)\left|\nabla u_{p}\right|^{p} \psi d x \mid \\
\leq & \left.\left|\int_{K}\left(u-\left(u_{p}, 0\right)\right)\right| \nabla u\right|^{p} \psi d x\left|+\left|\int_{K}\left(u_{p}, 0\right)\left(|\nabla u|^{p}-\left|\nabla u_{p}\right|^{p}\right) \psi d x\right| \rightarrow 0 .\right.
\end{aligned}
$$

Substituting this and (3.23) into (3.21), we see that the left hand side of (3.21) becomes

$$
\begin{aligned}
& \lim _{\varepsilon \rightarrow 0}\left[\int_{K} u \psi|\nabla u|^{p} d x+\frac{1}{\varepsilon^{p}} \int_{K} \psi\left(u u_{3}^{2}-u_{3} e_{3}\right) d x\right] \\
& =\int_{K}\left(u_{p}, 0\right)\left|\nabla u_{p}\right|^{p} \psi d x-\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon^{p}} \int_{K} \psi u_{3} e_{3} d x
\end{aligned}
$$

Comparing this with the right hand side of (3.21), we have

$$
\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon^{p}} \int_{K} \psi u_{3} e_{3} d x=0
$$

This is (1.4). Theorem 1.2 is proved.

## 4 A Preliminary Proposition

To present the convergent rate of $\left|u_{\varepsilon}^{\prime}\right| \rightarrow 1$ and $u_{\varepsilon 3} \rightarrow 0$ in $W^{1, p}$ sense when $\varepsilon \rightarrow 0$, we need the following
Proposition 4.1 Assume $u_{\varepsilon}$ is a minimizer of $E_{\varepsilon}(u, G)$ on $W$. If $E_{\varepsilon}\left(u_{\varepsilon}, K\right) \leq$ $C$ for some domain $K \subseteq G$. Then there exists a positive constant $C$ which is independent of $\varepsilon \in(0,1)$, such that

$$
\begin{equation*}
\frac{1}{p} \int_{K}\left|\nabla u_{\varepsilon}\right|^{p} d x+\frac{1}{\varepsilon^{p}} \int_{K} u_{\varepsilon 3}^{2} d x \leq C \varepsilon^{2 / p}+\frac{1}{p} \int_{K}\left|\nabla \frac{u_{\varepsilon}^{\prime}}{\left|u_{\varepsilon}^{\prime}\right|}\right|^{p} d x . \tag{4.1}
\end{equation*}
$$

Proof. Denote $w=\frac{u_{\varepsilon}^{\prime}}{\mid u_{\varepsilon}^{\prime}}$. Choose $R>0$ sufficiently small such that $\overline{B(x, 3 R)} \subset$ $K$. It follows from (3.6) that

$$
\begin{equation*}
\left|u_{\varepsilon}^{\prime}\right| \geq 1 / 2 \tag{4.2}
\end{equation*}
$$

on $B(x, 3 R)$ as $\varepsilon$ sufficiently small. This and (3.3) imply

$$
\begin{equation*}
\int_{B(x, 3 R)}|\nabla w|^{p} d x \leq 2^{p} \int_{B(x, 3 R)}\left|u_{\varepsilon}^{\prime}\right|^{p}|\nabla w|^{p} d x \leq C \int_{B(x, 3 R)}\left|\nabla u_{\varepsilon}\right|^{p} d x \leq C \tag{4.3}
\end{equation*}
$$

Applying (1.1) and the integral mean value theorem, we know that there is a constant $r \in(2 R, 3 R)$ such that

$$
\begin{equation*}
\frac{1}{p} \int_{\partial B(x, r)}\left|\nabla u_{\varepsilon}\right|^{p} d x+\frac{1}{2 \varepsilon^{p}} \int_{\partial B(x, r)} u_{\varepsilon 3}^{2} d x=C_{0}(r) E_{\varepsilon}\left(u_{\varepsilon}, B_{3 R} \backslash B_{2 R}\right) \leq C . \tag{4.4}
\end{equation*}
$$

Consider the functional

$$
E(\rho, B)=\frac{1}{p} \int_{B}\left(|\nabla \rho|^{2}+1\right)^{p / 2} d x+\frac{1}{2 \varepsilon^{p}} \int_{B}(1-\rho)^{2} d x,
$$

where $B=B(x, r)$. It is easy to prove that the minimizer $\rho_{1}$ of $E(\rho, B)$ on $W_{\left|u_{\varepsilon}^{\prime}\right|}^{1, p}\left(B, R^{+} \cup\{0\}\right)$ exists and solves

$$
\begin{align*}
-\operatorname{div}\left(v^{(p-2) / 2} \nabla \rho\right) & =\frac{1}{\varepsilon^{p}}(1-\rho) \quad \text { on } \quad B,  \tag{4.5}\\
\left.\rho\right|_{\partial B} & =\left|u_{\varepsilon}^{\prime}\right|, \tag{4.6}
\end{align*}
$$

where $v=|\nabla \rho|^{2}+1$. Since $1 / 2<\left|u_{\varepsilon}^{\prime}\right| \leq 1$, it follows from the maximum principle that on $\bar{B}$,

$$
\begin{equation*}
\frac{1}{2}<\rho_{1} \leq 1 \tag{4.7}
\end{equation*}
$$

Clearly, $\left(1-\left|u^{\prime}\right|\right)^{2} \leq\left(1-\left|u^{\prime}\right|^{2}\right)^{2}=u_{3}^{4} \leq u_{3}^{2}$. Thus, by noting that $\rho_{1}$ is a minimizer, and applying (1.1) we see easily that

$$
\begin{equation*}
E\left(\rho_{1}, B\right) \leq E\left(\left|u_{\varepsilon}^{\prime}\right|, B\right) \leq C E_{\varepsilon}\left(u_{\varepsilon}, B\right) \leq C . \tag{4.8}
\end{equation*}
$$

Multiplying (4.5) by $(\nu \cdot \nabla \rho)$, where $\rho$ denotes $\rho_{1}$, and integrating over $B$, we have

$$
\begin{align*}
-\int_{\partial B} v^{(p-2) / 2}(\nu \cdot \nabla \rho)^{2} d \xi & +\int_{B} v^{(n-2) / 2} \nabla \rho \cdot \nabla(\nu \cdot \nabla \rho) d x  \tag{4.9}\\
& =\frac{1}{\varepsilon^{p}} \int_{B}(1-\rho)(\nu \cdot \nabla \rho) d x,
\end{align*}
$$

where $\nu$ denotes the unit vector on $\bar{B}$, and it equals to the unit outside norm vector on $\partial B$.

Using (4.8) we obtain

$$
\begin{aligned}
& \left|\int_{B} v^{(p-2) / 2} \nabla \rho \cdot \nabla(\nu \cdot \nabla \rho) d x\right| \\
\leq & C \int_{B} v^{(n-2) / 2}|\nabla \rho|^{2} d x+\frac{1}{2}\left|\int_{B} v^{(p-2) / 2}(\nu \cdot \nabla v) d x\right| \\
\leq & C+\frac{1}{p}\left|\int_{B} \nu \cdot \nabla\left(v^{n / 2}\right) d x\right| \leq C+\frac{1}{p} \int_{B}\left|d i v\left(v^{p / 2} \nu\right)-v^{p / 2} d i v \nu\right| d x \\
\leq & C+\frac{1}{p} \int_{\partial B} v^{p / 2} d \xi .
\end{aligned}
$$

Combining (4.6), (4.4) and (4.8) we also have

$$
\begin{aligned}
& \left|\frac{1}{\varepsilon^{p}} \int_{B}(1-\rho)(\nu \cdot \nabla \rho) d x\right| \leq \frac{1}{2 \varepsilon^{p}}\left|\int_{B}(1-\rho)^{2} d i v \nu d x-\int_{\partial B}(1-\rho)^{2} d \xi\right| \\
\leq & \frac{1}{2 \varepsilon^{p}} \int_{B}(1-\rho)^{2}|d i v \nu| d x+\frac{1}{2 \varepsilon^{p}} \int_{\partial B}(1-\rho)^{2} d \xi \leq C .
\end{aligned}
$$

Substituting these into (4.9) yields

$$
\begin{equation*}
\left|\int_{\partial B} v^{(p-2) / 2}(\nu \cdot \nabla \rho)^{2} d \xi\right| \leq C+\frac{1}{p} \int_{\partial B} v^{p / 2} d \xi \tag{4.10}
\end{equation*}
$$

Applying (4.6), (4.4) and (4.10), we obtain for any $\delta \in(0,1)$,

$$
\begin{aligned}
& \int_{\partial B} v^{p / 2} d \xi=\int_{\partial B} v^{(p-2) / 2}\left[1+(\tau \cdot \nabla \rho)^{2}+(\nu \cdot \nabla \rho)^{2}\right] d \xi \\
& \leq \int_{\partial B} v^{(p-2) / 2} d \xi+\int_{\partial B} v^{(p-2) / 2}(\nu \cdot \nabla \rho)^{2} d \xi \\
& +\left(\int_{\partial B} v^{p / 2} d \xi\right)^{(p-2) / p}\left(\int_{\partial B}\left(\tau \cdot \nabla\left|u_{\varepsilon}^{\prime}\right|\right)^{p} d \xi\right)^{2 / p} \\
& \leq C(\delta)+\left(\frac{1}{p}+2 \delta\right) \int_{\partial B} v^{p / 2} d \xi
\end{aligned}
$$

where $\tau$ denotes the unit tangent vector on $\partial B$. Hence it follows by choosing $\delta>0$ so small that

$$
\begin{equation*}
\int_{\partial B} v^{p / 2} d \xi \leq C \tag{4.11}
\end{equation*}
$$

Now we multiply both sides of (4.5) by $(1-\rho)$ and integrate over $B$. Then

$$
\int_{B} v^{(p-2) / 2}|\nabla \rho|^{2} d x+\frac{1}{\varepsilon^{p}} \int_{B}(1-\rho)^{2} d x=-\int_{\partial B} v^{(p-2) / 2}(\nu \cdot \nabla \rho)(1-\rho) d \xi .
$$

From this, using (4.4), (4.6), (4.7) and (4.11) we obtain

$$
\begin{align*}
& E\left(\rho_{1}, B\right) \leq C\left|\int_{\partial B} v^{(p-2) / 2}(\nu \cdot \nabla \rho)(1-\rho) d \xi\right| \\
& \leq C\left|\int_{\partial B} v^{p / 2} d \xi\right|^{(p-1) / p}\left|\int_{\partial B}(1-\rho)^{2} d \xi\right|^{1 / p}  \tag{4.12}\\
& \leq C\left|\int_{\partial B}\left(1-\left|u_{\varepsilon}^{\prime}\right|\right)^{2} d \xi\right|^{1 / p} \leq C \varepsilon
\end{align*}
$$

Since $u_{\varepsilon}$ is a minimizer of $E_{\varepsilon}(u, G)$, we have

$$
E_{\varepsilon}\left(u_{\varepsilon}, G\right) \leq E_{\varepsilon}(U, G)
$$

where

$$
U=\left(\rho_{1} w, \sqrt{1-\rho_{1}^{2}}\right) \quad \text { on } \quad B ; \quad U=u_{\varepsilon} \quad \text { on } \quad G \backslash B .
$$

Namely,

$$
E_{\varepsilon}\left(u_{\varepsilon}, G\right) \leq E_{\varepsilon}\left(\rho_{1} w, B\right)+E_{\varepsilon}\left(u_{\varepsilon}, G \backslash B\right) .
$$

Hence

$$
\begin{align*}
& E_{\varepsilon}\left(u_{\varepsilon}, B\right) \leq E_{\varepsilon}\left(\rho_{1} w, B\right) \\
= & \frac{1}{p} \int_{B}\left(\left|\nabla \rho_{1}\right|^{2}+\rho_{1}^{2}|\nabla w|^{2}\right)^{p / 2} d x+\frac{1}{2 \varepsilon^{p}} \int_{B}\left(1-\rho_{1}^{2}\right) d x \tag{4.13}
\end{align*}
$$

where $w=\frac{u_{\varepsilon}^{\prime}}{\left|u_{\varepsilon}^{\prime}\right|}$. On one hand,

$$
\begin{align*}
& \int_{B}\left(\left|\nabla \rho_{1}\right|^{2}+\rho_{1}^{2}|\nabla w|^{2}\right)^{p / 2} d x-\int_{B}\left(\rho_{1}^{2}|\nabla w|^{2}\right)^{p / 2} d x \\
= & \frac{p}{2} \int_{B} \int_{0}^{1}\left(\left|\nabla \rho_{1}\right|^{2}+\rho_{1}^{2}|\nabla w|^{2}\right)^{(p-2) / 2} s  \tag{4.14}\\
& \left.+\left(\rho_{1}^{2}|\nabla w|^{2}\right)^{(p-2) / 2}(1-s)\right] d s\left|\nabla \rho_{1}\right|^{2} d x \\
\leq & C \int_{B}\left(\left|\nabla \rho_{1}\right|^{p}+\left|\nabla \rho_{1}\right|^{2}|\nabla w|^{p-2}\right) d x .
\end{align*}
$$

On the other hand, by using (4.12) and (4.3) we have

$$
\begin{equation*}
\int_{B}\left|\nabla \rho_{1}\right|^{2}|\nabla w|^{p-2} d x \leq\left(\int_{B}\left|\nabla \rho_{1}\right|^{p} d x\right)^{2 / p} \cdot\left(\int_{B}|\nabla w|^{p} d x\right)^{(p-2) / p} \leq C \varepsilon^{2 / p} \tag{4.15}
\end{equation*}
$$

Combining (4.13)-(4.15), we can derive

$$
E_{\varepsilon}\left(u_{\varepsilon}, B\right) \leq \frac{1}{p} \int_{B} \rho_{1}^{p}|\nabla w|^{p} d x+C \varepsilon^{2 / p} .
$$

Thus (4.1) can be seen by noticing (4.7).

## 5 Proof of Theorem 1.3

Assume that $u_{\varepsilon}$ is a minimizer, and $B=B(x, r)$. By noting $p>2$ and using Jensen's inequality, we have

$$
E_{\varepsilon}\left(u_{\varepsilon}, B\right) \geq \frac{1}{p} \int_{B}|\nabla h|^{p} d x+\frac{1}{p} \int_{B} h^{p}|\nabla w|^{p} d x+\frac{1}{p} \int_{B}\left|\nabla u_{3}\right|^{p} d x+\frac{1}{2 \varepsilon^{p}} \int_{B} u_{3}^{2} d x
$$

where $h=\left|u_{\varepsilon}^{\prime}\right|$. Thus, from (4.1) it follows that,

$$
\begin{align*}
& \frac{1}{p} \int_{B}\left(|\nabla h|^{p}+\left|\nabla u_{3}\right|^{p}\right) d x+\frac{1}{p} \int_{B}\left(h^{p}-1\right)|\nabla w|^{p} d x+\frac{1}{2 \varepsilon^{p}} \int_{B} u_{3}^{2} d x  \tag{5.1}\\
& \leq E_{\varepsilon}\left(u_{\varepsilon}, B\right)-\frac{1}{p} \int_{B}|\nabla w|^{p} d x \leq C \varepsilon^{2 / p} .
\end{align*}
$$

Since $\left|u_{3}\right| \leq 1$ and (1.1), we have

$$
\left|u_{3}(x)-u_{3}(y)\right| \leq C\left\|u_{3}\right\|_{W^{1, p}(K)}|x-y|^{1-2 / p} \leq C|x-y|^{1-2 / p}, \quad \forall x, y \in K .
$$

Hence, $u_{3}^{2}(x) \geq\left(\left|u_{3}(y)\right|-C \varepsilon^{1-2 / p}\right)^{2}$ when $x \in B(y, \varepsilon)$. Substituting this into (3.4) we obtain

$$
\pi\left(\left|u_{3}(y)\right|-C \varepsilon^{1-2 / p}\right)^{2} \varepsilon^{2} \leq \int_{B(y, \varepsilon)} u_{3}^{2}(x) d x \leq C \varepsilon^{p}
$$

for any $y \in K$. This implies

$$
\sup _{y \in K}\left|u_{3}(y)\right| \leq C \varepsilon^{1-2 / p} .
$$

Thus, by using (4.2) and (3.3), we have that for any constant $\delta \in(0,1)$,

$$
\begin{align*}
\frac{1}{p} \int_{B}\left(1-h^{p}\right)|\nabla w|^{p} d x & \leq \frac{2^{p}}{p} \int_{B}\left(1-h^{p}\right) h^{p}|\nabla w|^{p} d x  \tag{5.2}\\
& \leq C \int_{B} u_{3}^{2}\left|\nabla u_{\varepsilon}\right|^{p} d x \leq C \varepsilon^{1-2 / p}
\end{align*}
$$

Substituting this into (5.1), we can derive

$$
\begin{equation*}
\int_{B}|\nabla h|^{p} d x+\int_{B}\left|\nabla u_{3}\right|^{p} d x+\frac{1}{\varepsilon^{p}} \int_{B} u_{3}^{2} d x \leq C\left(\varepsilon^{1-2 / p}+\varepsilon^{2 / p}\right) . \tag{5.3}
\end{equation*}
$$

If $p \leq 4$, then we have finished. If $p>4$, we will prove
Theorem 5.1 Let $p_{0} \in(4,5)$ satisfy $p^{3}-4 p^{2}-2 p+4=0$. Then

$$
\begin{gathered}
\int_{B}|\nabla h|^{p} d x+\int_{B}\left|\nabla u_{3}\right|^{p} d x+\frac{1}{\varepsilon^{p}} \int_{B} u_{3}^{2} d x \leq C \varepsilon^{1-2 / p}, \quad \text { when } \quad p \in\left(4, p_{0}\right] ; \\
\int_{B}|\nabla h|^{p} d x+\int_{B}\left|\nabla u_{3}\right|^{p} d x+\frac{1}{\varepsilon^{p}} \int_{B} u_{3}^{2} d x \leq C \varepsilon^{\frac{2 p}{p^{2}-2}}, \quad \text { when } \quad p>p_{0} .
\end{gathered}
$$

Proof. Step 1. The idea of Proposition 4.1 is used. At first, from (5.3) it follows that

$$
\int_{B} u_{3}^{2} d x \leq C \varepsilon^{\frac{2}{p}+p}
$$

Using this and the integral mean theorem, we see that there exists $r_{2} \in(2 R, r)$ such that

$$
\int_{\partial B\left(x, r_{2}\right)} u_{3}^{2} d x \leq C \varepsilon^{\frac{2}{p}+p}
$$

Next, consider the minimizer $\rho_{2}$ of the functional

$$
E\left(\rho, B\left(x, r_{2}\right)\right)=\frac{1}{p} \int_{B\left(x, r_{2}\right)}\left(|\nabla \rho|^{2}+1\right)^{p / 2} d x+\frac{1}{2 \varepsilon^{p}} \int_{B\left(x, r_{2}\right)}(1-\rho)^{2} d x
$$

in $W_{\left|u_{\varepsilon}^{\prime}\right|}^{1, p}\left(B\left(x, r_{2}\right), R^{+} \cup\{0\}\right)$. By the same argument of (4.12) we also obtain

$$
E\left(\rho_{2}, B\left(x, r_{2}\right)\right) \leq C \varepsilon^{\frac{1}{p}\left(\frac{2}{p}+p\right)} .
$$

Then, similar to the derivation of (4.1) we can see that

$$
E_{\varepsilon}\left(u_{\varepsilon}, B\left(x, r_{2}\right)\right) \leq \frac{1}{p} \int_{B}|\nabla w|^{p} d x+C \varepsilon^{\frac{2}{p^{2}}\left(\frac{2}{p}+p\right)} .
$$

At last, by processing as the proof of (5.3) we have
$\int_{B\left(x, r_{2}\right)}|\nabla h|^{p} d x+\int_{B\left(x, r_{2}\right)}\left|\nabla u_{3}\right|^{p} d x+\frac{1}{\varepsilon^{p}} \int_{B\left(x, r_{2}\right)} u_{3}^{2} d x \leq C\left(\varepsilon^{1-2 / p}+\varepsilon^{\frac{2}{p^{2}}\left(\frac{2}{p}+p\right)}\right)$.
Step 2. Replacing (5.3) by the inequality above, and via the similar argument of Step 1, we also deduce that there exist $r_{j} \in\left(2 R, r_{j-1}\right)$ such that for any $j=1,2, \cdots$,

$$
\begin{equation*}
\int_{B\left(x, r_{j}\right)}|\nabla h|^{p} d x+\int_{B\left(x, r_{j}\right)}\left|\nabla u_{3}\right|^{p} d x+\frac{1}{\varepsilon^{p}} \int_{B\left(x, r_{j}\right)} u_{3}^{2} d x \leq C\left(\varepsilon^{1-2 / p}+\varepsilon^{a_{j}}\right) \tag{5.4}
\end{equation*}
$$

where $a_{1}=\frac{2}{p}$ and $a_{j}=\frac{2}{p^{2}}\left(a_{j-1}+p\right)$ for $j=2,3, \cdots$. Obviously, $\left\{a_{j}\right\}$ is a increasing and bounded sequence. So we see easily that its limit is $\frac{2 p}{p^{2}-2}$. Letting $j \rightarrow \infty$ in (5.4) we have proved Theorem 5.1.

Combining Theorem 5.1 and (5.3) yields that Theorem 1.3 is proved.

Acknowledgements. The research was supported by NSF (19271086) and Tianyuan Fund of Mathematics (A0324628)(China).

## References

[1] F.Bethuel, H.Brezis, F.Helein: Asymptotics for the minimization of a Ginzburg-Landau functional, Calc. Var. PDE., 1 (1993), 123-138.

EJQTDE, 2004 No. 16, p. 16
[2] F. Bethuel, H. Brezis, F. Helein: Ginzburg-Landau Vortices, Birkhauser. 1994.
[3] F. Bethuel, X. Zheng: Density of smooth functions between two manifolds in Sobolev space, J. Functional Anal., 80 (1988), 60-75.
[4] Y.Chen: Weak solutions to the evolution problem of harmonic maps, Math. Z., 201, (1989), 69-74.
[5] Y.Chen, M.Hong, N.Hungerbuhler: Heat flow of p-harmonic maps with values into spheres, Math. Z., 215 (1994), 25-35.
[6] Y.Chen, M.Struwe: Existence and partial regularity results for the heat flow for harmonic maps, Math. Z., 201, (1989), 83-103.
[7] F.Hang, F.Lin: Static theory for planar ferromagnets and antiferromagnets, Acta. Math. Sinica, English Series, 17, (2001), 541-580.
[8] M.Hong: Asymptotic behavior for minimizers of a Ginzburg-Landau type functional in higher dimensions associated with n-harmonic maps, Adv. in Diff. Eqns., 1 (1996), 611-634.
[9] S.Komineas, N.Papanicolaou: Vortex dynamics in two-dimensional antiferromagnets, Nonlinearity, 11, (1998), 265-290.
[10] Y.Lei: On the minimization of an energy functional J. Math. Anal. Appl. 293, (2004), 237-257.
[11] M.Misawa: Approximation of p-harmonic maps by the penalized equation, Nonlinear Anal. TMA., 47 (2001), 1069-1080.
[12] N.Papanicolaou, P.N.Spathis: Semitopological solutions in planar ferromagnets, Nonlinearity, 12, (1999), 285-302.
[13] M.Struwe: On the asymptotic behaviour of minimizers of the GinzburgLandau model in 2 dimensions, Diff. and Inter. Eqns., 7 (1994), 1613-1624.

