# Qualitative analysis of a mechanical system of coupled nonlinear oscillators 

Gheorghe Moroșanu ${ }^{1,2}$ and Cristian Vladimirescu ${ }^{\boxtimes 3}$<br>${ }^{1}$ Faculty of Mathematics and Computer Science, Babeș-Bolyai University, 1 M. Kogălniceanu Str., Cluj-Napoca, RO-400084, Romania<br>${ }^{2}$ Academy of Romanian Scientists, 3 Ilfov Str., Sector 5, Bucharest, RO-050045, Romania<br>${ }^{3}$ Department of Applied Mathematics, University of Craiova, 13 A.I. Cuza Str., Craiova, RO-200585, Romania

Received 1 December 2022, appeared 10 May 2023
Communicated by Bo Zhang


#### Abstract

In this paper we investigate nonlinear systems of second order ODEs describing the dynamics of two coupled nonlinear oscillators of a mechanical system. We obtain, under certain assumptions, some stability results for the null solution. Also, we show that in the presence of a time-dependent external force, every solution starting from sufficiently small initial data and its derivative are bounded or go to zero as the time tends to $+\infty$, provided that suitable conditions are satisfied. Our theoretical results are illustrated with numerical simulations.


Keywords: coupled oscillators, uniform stability, asymptotic stability, uniform asymptotic stability.
2020 Mathematics Subject Classification: 34C15, 34D20.

## 1. Introduction

Consider a mechanical system of coupled nonlinear oscillators, as shown in Figure 1.1. Specifically, the block of mass $m_{1}$ is anchored to a fixed horizontal wall and the block of mass $m_{2}$ by springs and dampers, and the block of mass $m_{2}$ is also attached to the wall by a pair of springs and dampers. Suppose that the stiffnesses and the dampings are represented by the functions $k_{i}: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}$and $d_{i}: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}, i \in\{1,2,3\}$, and $\widehat{g}_{i}: \mathbb{R}_{+} \times \mathbb{R} \times \mathbb{R} \rightarrow \mathbb{R}, i \in\{1,2\}$, denote external forces acting on the blocks. One may also consider an external force $\widehat{f}(t)$ acting on the block of mass $m_{1}$, but for the moment, we restrict our attention to the case $\widehat{f} \equiv 0$. We assume that when the two blocks are in their equilibrium positions, the springs and the dampers are also in their equilibrium positions. Let $x(t)$ and $y(t)$ be the vertical displacements of the blocks from their equilibrium positions.

[^0]

Figure 1.1: A mechanical system of coupled nonlinear oscillators

Then the system of ODEs describing the motion is (see, e.g., [27])

$$
\left\{\begin{array}{l}
m_{1} \ddot{x}+k_{1}(t) x+2 d_{1}(t) \dot{x}-k_{3}(t)(y-x)-2 d_{3}(t)(\dot{y}-\dot{x})=\widehat{g}_{1}(t, x, y), \\
m_{2} \ddot{y}+2 k_{2}(t) y+4 d_{2}(t) \dot{y}+k_{3}(t)(y-x)+2 d_{3}(t)(\dot{y}-\dot{x})=\widehat{g}_{2}(t, x, y),
\end{array}\right.
$$

or

$$
\left\{\begin{array}{l}
\ddot{x}+2 f_{1}(t) \dot{x}-f_{3}(t) \dot{y}+\beta(t) x-\gamma_{1}(t) y+g_{1}(t, x, y)=0,  \tag{1.1}\\
\ddot{y}+2 f_{2}(t) \dot{y}-f_{4}(t) \dot{x}-\gamma_{2}(t) x+\delta(t) y+g_{2}(t, x, y)=0,
\end{array}\right.
$$

where

$$
\begin{aligned}
f_{1}(t) & :=\frac{1}{m_{1}}\left(d_{1}(t)+d_{3}(t)\right), & f_{2}(t) & :=\frac{1}{m_{2}}\left(2 d_{2}(t)+d_{3}(t)\right), \\
f_{3}(t) & :=\frac{2}{m_{1}} d_{3}(t), & f_{4}(t) & :=\frac{2}{m_{2}} d_{3}(t), \\
\beta(t) & :=\frac{1}{m_{1}}\left(k_{1}(t)+k_{3}(t)\right), & \delta(t) & :=\frac{1}{m_{2}}\left(2 k_{2}(t)+k_{3}(t)\right), \\
\gamma_{1}(t) & :=\frac{1}{m_{1}} k_{3}(t), & \gamma_{2}(t) & :=\frac{1}{m_{2}} k_{3}(t), \\
g_{1}(t, x, y) & :=-\frac{1}{m_{1}} \widehat{g}_{1}(t, x, y), & g_{2}(t, x, y) & :=-\frac{1}{m_{2}} \widehat{g}_{2}(t, x, y) .
\end{aligned}
$$

The general case of a single 1-D damped nonlinear oscillator is described by the following equation which is well-known in the literature

$$
\begin{equation*}
\ddot{x}+2 f^{*}(t) \dot{x}+\beta^{*}(t) x+g^{*}(t, x)=0, \quad t \in \mathbb{R}_{+} . \tag{1.2}
\end{equation*}
$$

T. A. Burton and T. Furumochi [2] introduced a new method, based on the Schauder fixed point theorem, to study the stability of the null solution of Eq. (1.2) in the case $\beta^{*}(t)=1$. In [14] we reported new stability results for the same equation. Our approach was based on elementary arguments only, involving in particular some Bernoulli type differential inequalities. In [15] we considered Eq. (1.2) under more general assumptions, which required more
sophisticated arguments. For other investigations regarding the asymptotic stability of the equilibrium of a single damped nonlinear oscillator, we refer the reader to $[7,8,10,11,24]$, and the references therein.

In the present paper, in Section 2 we will study the stability of the null solution of system (1.1), by two approaches, based on classical differential inequalities and on Lyapunov's method. For other results regarding the asymptotic stability of the equilibria of coupled damped nonlinear oscillators, we refer the reader to [9,16,17,20-23,25], and the references therein. For fundamental concepts and results in stability theory we refer the reader to [1,3,5,6,13,19].

In Section 3 we will consider that the block of mass $m_{1}$ is subject to the action of a time dependent external force $\widehat{f}: \mathbb{R}_{+} \rightarrow \mathbb{R}$. In this case, the system of ODEs describing the dynamics of the mechanical system is

$$
\left\{\begin{array}{l}
\ddot{x}+2 f_{1}(t) \dot{x}-f_{3}(t) \dot{y}+\beta(t) x-\gamma_{1}(t) y-f(t)+g_{1}(t, x, y)=0,  \tag{1.3}\\
\ddot{y}+2 f_{2}(t) \dot{y}-f_{4}(t) \dot{x}-\gamma_{2}(t) x+\delta(t) y+g_{2}(t, x, y)=0,
\end{array}\right.
$$

with the same functions as before, and $f(t):=\frac{1}{m_{1}} \widehat{f}(t)$, and we will derive certain qualitative properties of the solutions of system (1.3) with initial data small enough.

The model in Figure 1.1 could be used, e.g., to describe the dynamics in vertical direction of vibration reduction systems for horizontal cranes with loadings suspended in two sides $[12,28]$. For other models of coupled oscillators or for models from electric circuit theory, structural dynamics, described by systems of type (1.1) or (1.3), we refer the reader to the monographs [4,18,26].

## 2. A stability result for the system (1.1)

In this section we shall use the following hypotheses.
(H1) $f_{i} \in C^{1}\left(\mathbb{R}_{+}\right), f_{j} \in C\left(\mathbb{R}_{+}\right), f_{i}(t) \geq 0, f_{j}(t) \geq 0, \forall t \in \mathbb{R}_{+}$, and $\int_{0}^{+\infty} f_{j}(t) \mathrm{d} t<+\infty$, $\forall i \in\{1,2\}, \forall j \in\{3,4\}$;
(H2) there exist constants $h, K_{1}, K_{2} \geq 0$ such that

$$
\left|\dot{f}_{i}(t)+f_{i}^{2}(t)\right| \leq K_{i} \tilde{f}(t), \forall t \in[h,+\infty), \forall i \in\{1,2\},
$$

where $\widetilde{f}(t):=\min \left\{f_{1}(t), f_{2}(t)\right\}, \forall t \in \mathbb{R}_{+} ;$
(H3) $\int_{0}^{+\infty} \widetilde{f}(t) \mathrm{d} t=+\infty$.
(H4) $\beta, \delta \in C^{1}\left(\mathbb{R}_{+}\right), \beta, \delta$ are decreasing and

$$
\beta(t) \geq \beta_{0}>0, \delta(t) \geq \delta_{0}>0, \quad \forall t \in \mathbb{R}_{+},
$$

where $\beta_{0}, \delta_{0}$ are constants such that

$$
\frac{K_{1}}{\sqrt{\beta_{0}}}+\frac{K_{2}}{\sqrt{\delta_{0}}}<1 ;
$$

(H5) $\gamma_{i} \in C\left(\mathbb{R}_{+}\right), \gamma_{i}(t) \geq 0, \forall t \in \mathbb{R}_{+}$, and $\int_{0}^{+\infty} \gamma_{i}(t) \mathrm{d} t<+\infty, \forall i \in\{1,2\}$;
(H6) $g_{i}=g_{i}(t, x, y) \in C\left(\mathbb{R}_{+} \times \mathbb{R} \times \mathbb{R}\right), g_{i}$ are locally Lipschitzian with respect to $x, y, i \in$ $\{1,2\}$, and fulfill the relations

$$
\begin{align*}
& \left|g_{1}(t, x, y)\right| \leq r_{1}(t) \mathrm{O}(|x|), \quad \forall t \in \mathbb{R}_{+}, \quad \forall y \in \mathbb{R},  \tag{2.1}\\
& \left|g_{2}(t, x, y)\right| \leq r_{2}(t) \mathrm{O}(|y|), \quad \forall t \in \mathbb{R}_{+}, \forall x \in \mathbb{R}, \tag{2.2}
\end{align*}
$$

where $r_{i} \in C\left(\mathbb{R}_{+}\right), r_{i}(t) \geq 0, \forall t \in \mathbb{R}_{+}, \int_{0}^{+\infty} r_{i}(t) \mathrm{d} t<+\infty, \forall i \in\{1,2\}$, and $\mathrm{O}(|x|)$ denotes the big-O Landau symbol as $x \rightarrow 0$ (similarly for $\mathrm{O}(|y|)$ );
(H7) There is a $p>0$, such that $f_{i}(t) \geq p, \forall t \geq 0, \forall i \in\{1,2\}$.
Remark 2.1. If (H1) and (H2) hold, then $f_{i}, \dot{f}_{i}$ are bounded, $i \in\{1,2\}$. Indeed, by (H2) we see that

$$
\left(t \geq h, f_{i}(t)>K_{i}\right) \Longrightarrow \dot{f}_{i}(t)<0
$$

This, combined with (H1), implies

$$
f_{i}(t) \leq M_{i}:=\max \left\{f_{i}(h), K_{i}\right\}, \quad \forall t \geq h .
$$

So, using again (H2), we obtain

$$
\left|\dot{f}_{i}(t)\right| \leq 2 M_{i}^{2}, \quad \forall t \geq h
$$

This concludes the proof, since, by (H1), $f_{i}, \dot{f}_{i} \in C[0, h], i \in\{1,2\}$.
Remark 2.2. Since we are going to discuss the stability of the null solution of system (1.1) and the large-time behavior of the solutions to (1.3) starting from small initial data, we can replace the inequalities (2.1) and (2.2) by

$$
\begin{equation*}
\left|g_{1}(t, x, y)\right| \leq r_{1}(t)|x|, \quad\left|g_{2}(t, x, y)\right| \leq r_{2}(t)|y|, \quad \forall t \in \mathbb{R}_{+}, \forall x, y \in \mathbb{R}, \tag{2.3}
\end{equation*}
$$

possibly with $M_{i} r_{i}(t)$ instead of $r_{i}(t)$, where $M_{i}>0$, and some functions $\widetilde{g}_{i}$ instead of $g_{i}$, $\forall i \in\{1,2\}$.

Indeed, from (2.1) there exist $M_{1}, a_{1}>0$, such that

$$
\left|g_{1}(t, x, y)\right| \leq r_{1}(t) M_{1}|x|, \quad \text { if }|x|<a_{1} .
$$

If we define the function $\widetilde{g}_{1}: \mathbb{R}_{+} \times \mathbb{R} \times \mathbb{R} \rightarrow \mathbb{R}$ as

$$
\widetilde{g}_{1}(t, x, y):=\left\{\begin{array}{l}
g_{1}\left(t, a_{1}, y\right), \text { if } x \geq a_{1} \\
g_{1}(t, x, y), \text { if }|x|<a_{1} \\
g_{1}\left(t,-a_{1}, y\right), \text { if } x \leq-a_{1}
\end{array}\right.
$$

for all $t \geq 0, y \in \mathbb{R}$, then

$$
\left|\widetilde{g}_{1}(t, x, y)\right| \leq r_{1}(t) M_{1}|x|, \quad \forall(t, x, y) \in \mathbb{R}_{+} \times \mathbb{R} \times \mathbb{R},
$$

$\widetilde{g}_{1} \in C\left(\mathbb{R}_{+} \times \mathbb{R} \times \mathbb{R}\right)$, and $\widetilde{g}_{1}$ is locally Lipschitzian in $x, y$. Similar reasonings work for the functions $g_{2}$ and $r_{2}$, possibly with another constant $a_{2}$.

### 2.1. A stability result via differential inequalities

We can state and prove the following stability result.

## Theorem 2.3.

a) Suppose that the hypotheses (H1), (H2), (H4)-(H6) are satisfied. Then the null solution of the system (1.1) is uniformly stable.
b) If the hypotheses (H1)-(H6) are fulfilled, then the null solution of (1.1) is asymptotically stable.
c) If the hypotheses (H1), (H2), (H4)-(H7) are fulfilled, then the null solution of (1.1) is uniformly asymptotically stable.
Proof. By using the following transformation (inspired from [2])

$$
\left\{\begin{array}{l}
\dot{x}=u-f_{1}(t) x  \tag{2.4}\\
\dot{u}=\left[\dot{f}_{1}(t)+f_{1}^{2}(t)-\beta(t)\right] x-f_{1}(t) u+\left[\gamma_{1}(t)-f_{2}(t) f_{3}(t)\right] y+f_{3}(t) v-g_{1}(t, x, y) \\
\dot{y}=v-f_{2}(t) y \\
\dot{v}=\left[\gamma_{2}(t)-f_{1}(t) f_{4}(t)\right] x+f_{4}(t) u+\left[\dot{f}_{2}(t)+f_{2}^{2}(t)-\delta(t)\right] y-f_{2}(t) v-g_{2}(t, x, y)
\end{array}\right.
$$

the system (1.1) becomes

$$
\begin{equation*}
\dot{z}=A(t) z+B(t) z+F(t, z) \tag{2.5}
\end{equation*}
$$

where

$$
\begin{gathered}
z=\left(\begin{array}{l}
x \\
u \\
y \\
v
\end{array}\right), \quad A(t)=\left(\begin{array}{cccc}
-f_{1}(t) & 1 & 0 & 0 \\
-\beta(t) & -f_{1}(t) & \gamma_{1}(t) & 0 \\
0 & 0 & -f_{2}(t) & 1 \\
\gamma_{2}(t) & 0 & -\delta(t) & -f_{2}(t)
\end{array}\right), \\
B(t)=\left(\begin{array}{ccc}
0 & 0 & 0 \\
f_{1}(t)+f_{1}^{2}(t) & 0 & -f_{2}(t) f_{3}(t) \\
0 & 0 & f_{3}(t) \\
-f_{1}(t) f_{4}(t) & f_{4}(t) & \dot{f}_{2}(t)+f_{2}^{2}(t) \\
0
\end{array}\right), \quad F(t, z)=\left(\begin{array}{c}
0 \\
-g_{1}(t, x, y) \\
0 \\
-g_{2}(t, x, y)
\end{array}\right) .
\end{gathered}
$$

Using the boundedness of the functions $f_{i}, \dot{f_{i}}, f_{j}, \beta, \gamma_{i}, \delta, r_{i}, \forall i \in\{1,2\}, \forall j \in\{3,4\}$, we easily deduce that our stability question of the null solution of the system (1.1) reduces to the stability of the null solution $z(t)=0$ of the system (2.5) .

Let $t_{0} \geq 0$ and

$$
\mathrm{Z}\left(t, t_{0}\right)=\left(a_{i j}\left(t, t_{0}\right)\right)_{i, j \in \overline{1,4^{\prime}}} t \geq t_{0}
$$

be the fundamental matrix of the system

$$
\begin{equation*}
\dot{z}=A(t) z, \tag{2.6}
\end{equation*}
$$

which equals the identity matrix for $t=t_{0}$. Then we deduce

$$
\begin{align*}
& \beta(t) a_{11}^{2}\left(t, t_{0}\right)+a_{21}^{2}\left(t, t_{0}\right)+\delta(t) a_{31}^{2}\left(t, t_{0}\right)+a_{41}^{2}\left(t, t_{0}\right) \leq \beta\left(t_{0}\right) \mathrm{e}^{\int_{t_{0}}^{t}\left[-2 \tilde{f}(u)+\frac{\gamma(u)}{\sqrt{\zeta}(u)}\right] \mathrm{d} u,}  \tag{2.7}\\
& \left.\beta(t) a_{12}^{2}\left(t, t_{0}\right)+a_{22}^{2}\left(t, t_{0}\right)+\delta(t) a_{32}^{2}\left(t, t_{0}\right)+a_{42}^{2}\left(t, t_{0}\right) \leq \mathrm{e}^{\int_{t_{0}}^{t}\left[-2 \tilde{f}(u)+\frac{\gamma(u)}{\sqrt{\zeta}(u)}\right]}\right] \mathrm{d} u,  \tag{2.8}\\
& \beta(t) a_{13}^{2}\left(t, t_{0}\right)+a_{23}^{2}\left(t, t_{0}\right)+\delta(t) a_{33}^{2}\left(t, t_{0}\right)+a_{43}^{2}\left(t, t_{0}\right) \leq \delta\left(t_{0}\right) \mathrm{e}^{\int_{t_{0}}^{t}\left[-2 \tilde{f}(u)+\frac{\gamma(u)}{\sqrt{\zeta(u u)}}\right] \mathrm{d} u},  \tag{2.9}\\
& \beta(t) a_{14}^{2}\left(t, t_{0}\right)+a_{24}^{2}\left(t, t_{0}\right)+\delta(t) a_{34}^{2}\left(t, t_{0}\right)+a_{44}^{2}\left(t, t_{0}\right) \leq \mathrm{e}^{\int_{t_{0}}^{t}\left[-2 \tilde{f}(u)+\frac{\gamma(u)}{\sqrt{\zeta}(u)}\right]} \mathrm{d} u, \tag{2.10}
\end{align*}
$$

for all $t \geq t_{0}$, where $\gamma(t):=\max \left\{\gamma_{1}(t), \gamma_{2}(t)\right\}, \zeta(t):=\min \{\beta(t), \delta(t)\}, \forall t \in \mathbb{R}_{+}$.
Indeed, from (2.6) we get the following system

$$
\left\{\begin{array}{l}
\dot{a}_{11}\left(t, t_{0}\right)=-f_{1}(t) a_{11}\left(t, t_{0}\right)+a_{21}\left(t, t_{0}\right)  \tag{2.11}\\
\dot{a}_{21}\left(t, t_{0}\right)=-\beta(t) a_{11}\left(t, t_{0}\right)-f_{1}(t) a_{21}\left(t, t_{0}\right)+\gamma_{1}(t) a_{31}\left(t, t_{0}\right) \\
\dot{a}_{31}\left(t, t_{0}\right)=-f_{2}(t) a_{31}\left(t, t_{0}\right)+a_{41}\left(t, t_{0}\right) \\
\dot{a}_{41}\left(t, t_{0}\right)=\gamma_{2}(t) a_{11}\left(t, t_{0}\right)-\delta(t) a_{31}\left(t, t_{0}\right)-f_{2}(t) a_{41}\left(t, t_{0}\right) .
\end{array}\right.
$$

From the first two equations of (2.11) and hypothesis (H4) we get

$$
\begin{align*}
& \frac{1}{2} \frac{\mathrm{~d}}{\mathrm{~d} t}\left[\beta(t) a_{11}^{2}\left(t, t_{0}\right)+a_{21}^{2}\left(t, t_{0}\right)\right] \\
& \quad \leq-f_{1}(t)\left[\beta(t) a_{11}^{2}\left(t, t_{0}\right)+a_{21}^{2}\left(t, t_{0}\right)\right]+\gamma_{1}(t) a_{21}\left(t, t_{0}\right) a_{31}\left(t, t_{0}\right) \tag{2.12}
\end{align*}
$$

and, similarly,

$$
\begin{align*}
& \frac{1}{2} \frac{\mathrm{~d}}{\mathrm{~d} t}\left[\delta(t) a_{31}^{2}\left(t, t_{0}\right)+a_{41}^{2}\left(t, t_{0}\right)\right] \\
& \quad \leq-f_{2}(t)\left[\delta(t) a_{31}^{2}\left(t, t_{0}\right)+a_{41}^{2}\left(t, t_{0}\right)\right]+\gamma_{2}(t) a_{11}\left(t, t_{0}\right) a_{41}\left(t, t_{0}\right) . \tag{2.13}
\end{align*}
$$

By relations (2.12) and (2.13) we obtain successively

$$
\begin{aligned}
\frac{1}{2} \frac{\mathrm{~d}}{\mathrm{~d} t} & {\left[\beta(t) a_{11}^{2}\left(t, t_{0}\right)+a_{21}^{2}\left(t, t_{0}\right)+\delta(t) a_{31}^{2}\left(t, t_{0}\right)+a_{41}^{2}\left(t, t_{0}\right)\right] } \\
\leq & -f_{1}(t)\left[\beta(t) a_{11}^{2}\left(t, t_{0}\right)+a_{21}^{2}\left(t, t_{0}\right)\right]-f_{2}(t)\left[\delta(t) a_{31}^{2}\left(t, t_{0}\right)+a_{41}^{2}\left(t, t_{0}\right)\right] \\
& +\gamma_{1}(t) a_{21}\left(t, t_{0}\right) a_{31}\left(t, t_{0}\right)+\gamma_{2}(t) a_{11}\left(t, t_{0}\right) a_{41}\left(t, t_{0}\right) \\
\leq & {\left[-\widetilde{f}(t)+\frac{\gamma(t)}{2 \sqrt{\zeta(t)}}\right]\left[\beta(t) a_{11}^{2}\left(t, t_{0}\right)+a_{21}^{2}\left(t, t_{0}\right)+\delta(t) a_{31}^{2}\left(t, t_{0}\right)+a_{41}^{2}\left(t, t_{0}\right)\right] }
\end{aligned}
$$

for all $t \geq t_{0}$, and (2.7) follows immediately. The inequalities (2.8)-(2.10) can be derived in the same way.

Let $\|\cdot\|_{0}$ be the norm in $\mathbb{R}^{4}$ defined by

$$
\begin{equation*}
\|z\|_{0}=\left(\beta_{0} x^{2}+u^{2}+\delta_{0} y^{2}+v^{2}\right)^{1 / 2}, \quad \text { for } z=(x, u, y, v)^{\top}, \tag{2.14}
\end{equation*}
$$

which is equivalent to the Euclidean norm.
For $z_{0}=\left(x_{0}, u_{0}, y_{0}, v_{0}\right)^{\top} \in \mathbb{R}^{4}$, from (2.7)-(2.10) and (H4), we deduce

$$
\begin{equation*}
\left\|Z\left(t, t_{0}\right) z_{0}\right\|_{0} \leq \lambda\left\|z_{0}\right\|_{0} \sqrt{\beta\left(t_{0}\right)+\delta\left(t_{0}\right)+2} \mathrm{e}^{\int_{t_{0}}^{t}}\left[-\tilde{f}(u)+\frac{\gamma(u)}{2 \sqrt{\xi(u)}}\right] \mathrm{d} u, \quad \forall t \geq t_{0} \tag{2.15}
\end{equation*}
$$

where $\lambda:=\max \left\{1,1 / \sqrt{\beta_{0}}, 1 / \sqrt{\delta_{0}}\right\}$,

$$
\begin{align*}
& \left\|Z\left(t, t_{0}\right) Z\left(s, t_{0}\right)^{-1} e_{2}\right\|_{0} \leq \mathrm{e}^{\int_{s}^{t}\left[-\tilde{f}(u)+\frac{\gamma(u)}{2 \sqrt{\xi(u)}}\right] \mathrm{d} u},  \tag{2.16}\\
& \left\|Z\left(t, t_{0}\right) Z\left(s, t_{0}\right)^{-1} e_{4}\right\|_{0} \leq \mathrm{e}^{\int_{s}^{t}\left[-\tilde{f}(u)+\frac{\gamma(u)}{2 \sqrt{\xi(u)}}\right] \mathrm{d} u},
\end{align*}
$$

for all $t \geq s \geq t_{0} \geq 0$, where $e_{2}=(0,1,0,0)^{\top}, e_{4}=(0,0,0,1)^{\top}$.

Proof of $a$ ). Let $z_{0} \neq 0$ with $\left\|z_{0}\right\|_{0}$ small enough, $t_{0} \geq 0$, and $z\left(t, t_{0}, z_{0}\right)=\left(x\left(t, t_{0}, z_{0}\right), u\left(t, t_{0}, z_{0}\right)\right.$, $\left.y\left(t, t_{0}, z_{0}\right), v\left(t, t_{0}, z_{0}\right)\right)^{\top}$ be the unique solution of (2.5) which equals $z_{0}$ for $t=t_{0}$.

From the continuity and the boundedness of the functions $f_{i}, \dot{f}_{i}, f_{j}, \beta, \gamma_{i}, \delta, r_{i}, \forall i \in\{1,2\}$, $\forall j \in\{3,4\}$, there exists $\psi: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}$a continuous and bounded function, such that

$$
\|A(t) z+B(t) z+F(t, z)\|_{0} \leq \psi(t)\|z\|_{0}, \quad \forall(t, z) \in \mathbb{R}_{+} \times \mathbb{R}^{4}
$$

By applying a classical result of global existence in the future to system (2.5) (see, e.g., [3, Corollary, p. 53]) it follows that $z\left(t, t_{0}, z_{0}\right)$ exists on the whole interval $\left[t_{0},+\infty\right)$.

We have

$$
\begin{equation*}
z\left(t, t_{0}, z_{0}\right)=Z\left(t, t_{0}\right) z_{0}+\int_{t_{0}}^{t} Z\left(t, t_{0}\right) Z\left(s, t_{0}\right)^{-1}\left[B(s) z\left(s, t_{0}, z_{0}\right)+F\left(s, z\left(s, t_{0}, z_{0}\right)\right)\right] \mathrm{d} s, \tag{2.17}
\end{equation*}
$$

for all $t \geq t_{0}$.
From the relations (2.15)-(2.17) we get

$$
\begin{align*}
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0} \leq & \left.\lambda\left\|z_{0}\right\|_{0} \sqrt{\beta\left(t_{0}\right)+\delta\left(t_{0}\right)+2} \mathrm{e}^{\int_{t_{0}}^{t}\left[-\tilde{f}(u)+\frac{\gamma(u)}{2 \sqrt{\zeta(u)}}\right]}\right] \mathrm{d} u \\
& \times \int_{t_{0}}^{t} \mathrm{e}^{\int_{s}^{t}\left[-\tilde{f}(u)+\frac{\gamma(u)}{2 \sqrt{\zeta(u)}}\right]} \mathrm{d} u \\
& +f_{1}(s) f_{4}(s)\left|x\left(s, t_{0}, z_{0}\right)\right|+f_{2}^{2}(s) f_{3}(s)\left|y\left(s, t_{0}, z_{0}\right)\right| \\
& +f_{3}(s)\left|v\left(s, t_{0}, z_{0}\right)\right|+f_{4}(s)\left|u\left(s, t_{0}, z_{0}\right)\right| \\
& +\left|g_{1}\left(s, x\left(s, t_{0}, z_{0}\right), y\left(s, t_{0}, z_{0}\right)\right)\right| \\
& \left.+\left|g_{2}\left(s, x\left(s, t_{0}, z_{0}\right), y\left(s, t_{0}, z_{0}\right)\right)\right|\right] \mathrm{d} s \tag{2.18}
\end{align*}
$$

for all $t \geq t_{0}$.
In what follows we consider two cases.
Case 1: $0 \leq t_{0}<h$. Since $f_{i} \in C^{1}\left[t_{0}, h\right], f_{j}, \beta, \gamma_{i}, \delta, \in C\left[t_{0}, h\right], g_{i} \in C\left(\left[t_{0}, h\right] \times \mathbb{R} \times \mathbb{R}\right), \forall i \in\{1,2\}$, $\forall j \in\{3,4\}$, from (2.18) it results that

$$
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0} \leq \lambda D_{1} \sqrt{\beta\left(t_{0}\right)+\delta\left(t_{0}\right)+2}\left\|z_{0}\right\|_{0}+D \int_{t_{0}}^{t}\left\|z\left(s, t_{0}, z_{0}\right)\right\|_{0} \mathrm{~d} s, \quad \forall t \in\left[t_{0}, h\right]
$$

with $D, D_{1}$ positive constants. Using the Gronwall lemma we get

$$
\begin{equation*}
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0} \leq \lambda D_{1} \sqrt{\beta\left(t_{0}\right)+\delta\left(t_{0}\right)+2}\left\|z_{0}\right\|_{0} \mathrm{e}^{D h}, \quad \forall t \in\left[t_{0}, h\right] . \tag{2.19}
\end{equation*}
$$

For all $t \geq h$, from the relation (2.18) and the hypothesis (H2) we deduce

$$
\begin{align*}
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0} \leq & \lambda \sqrt{\beta(h)+\delta(h)+2}\left\|z\left(h, t_{0}, z_{0}\right)\right\|_{0} \mathrm{e}^{\int_{h}^{t}\left[-\tilde{f}(u)+\frac{\gamma(u)}{2 \sqrt{\zeta(u)}}\right] \mathrm{d} u} \\
& +\int_{h}^{t} \mathrm{e}^{\int_{s}^{t}\left[-\widetilde{f}(u)+\frac{\gamma(u)}{2 \sqrt{\zeta(u)}}\right] \mathrm{d} u}\left[K_{1} \widetilde{f}(s)\left|x\left(s, t_{0}, z_{0}\right)\right|+K_{2} \widetilde{f}(s)\left|y\left(s, t_{0}, z_{0}\right)\right|\right. \\
& +f_{1}(s) f_{4}(s)\left|x\left(s, t_{0}, z_{0}\right)\right|+f_{2}(s) f_{3}(s)\left|y\left(s, t_{0}, z_{0}\right)\right| \\
& +f_{3}(s)\left|v\left(s, t_{0}, z_{0}\right)\right|+f_{4}(s)\left|u\left(s, t_{0}, z_{0}\right)\right| \\
& +\left|g_{1}\left(s, x\left(s, t_{0}, z_{0}\right), y\left(s, t_{0}, z_{0}\right)\right)\right| \\
& \left.+\left|g_{2}\left(s, x\left(s, t_{0}, z_{0}\right), y\left(s, t_{0}, z_{0}\right)\right)\right|\right] \mathrm{d} s . \tag{2.20}
\end{align*}
$$

By (2.3) and (2.20) we obtain

$$
\begin{align*}
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0} \leq & \lambda \sqrt{\beta(h)+\delta(h)+2}\left\|z\left(h, t_{0}, z_{0}\right)\right\|_{0} \mathrm{e}^{f_{h}^{t}}\left[-\tilde{f}(u)+\frac{\gamma(u)}{2 \sqrt{\xi(u)}}\right] \mathrm{d} u \\
& +\int_{h}^{t} \mathrm{e}^{\int_{s}^{t}\left[-\widetilde{f}(u)+\frac{\gamma(u)}{2 \sqrt{5(u)}}\right] \mathrm{d} u}\left[\left(\frac{K_{1}}{\sqrt{\beta_{0}}}+\frac{K_{2}}{\sqrt{\delta_{0}}}\right) \widetilde{f}(s)\right. \\
& +\frac{f_{1}(s) f_{4}(s)}{\sqrt{\beta_{0}}}+\frac{f_{2}(s) f_{3}(s)}{\sqrt{\delta_{0}}}+f_{3}(s)+f_{4}(s) \\
& \left.+\frac{r_{1}(s)}{\sqrt{\beta_{0}}}+\frac{r_{2}(s)}{\sqrt{\delta_{0}}}\right]\left\|z\left(s, t_{0}, z_{0}\right)\right\|_{0} \mathrm{~d} s \\
= & : \sigma(t), \quad \forall t \geq h . \tag{2.21}
\end{align*}
$$

Straightforward calculations lead us to

$$
\begin{gather*}
\dot{\sigma}(t) \leq \omega(t) \sigma(t), \quad \forall t \geq h,  \tag{2.22}\\
\sigma(h)=\lambda \sqrt{\beta(h)+\delta(h)+2}\left\|z\left(h, t_{0}, z_{0}\right)\right\|_{0} .
\end{gather*}
$$

where

$$
\begin{aligned}
\omega(t) & :=-K \widetilde{f}(t)+\varphi(t), \quad \forall t \geq 0, K=1-\frac{K_{1}}{\sqrt{\beta_{0}}}-\frac{K_{2}}{\sqrt{\delta_{0}}} \\
\varphi(t) & :=\frac{\gamma(t)}{2 \sqrt{\zeta(t)}}+\frac{f_{1}(t) f_{4}(t)}{\sqrt{\beta_{0}}}+\frac{f_{2}(t) f_{3}(t)}{\sqrt{\delta_{0}}}+f_{3}(t)+f_{4}(t)+\frac{r_{1}(t)}{\sqrt{\beta_{0}}}+\frac{r_{2}(t)}{\sqrt{\delta_{0}}}, \quad \forall t \geq 0
\end{aligned}
$$

From (2.21) and (2.22) using classical differential inequalities, we obtain

$$
\begin{equation*}
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0} \leq \lambda \sqrt{\beta(h)+\delta(h)+2}\left\|z\left(h, t_{0}, z_{0}\right)\right\|_{0} \mathrm{e}^{-K \int_{h}^{t} \tilde{f}(s) \mathrm{ds}} \mathrm{e}^{\int_{h}^{t} \varphi(s) \mathrm{d} s}, \quad \forall t \geq h \tag{2.23}
\end{equation*}
$$

It is readily seen from the hypotheses (H1), (H5), (H6), and Remark 2.1, that

$$
\int_{h}^{+\infty} \varphi(s) \mathrm{d} s<+\infty .
$$

Let $\varepsilon>0$ be arbitrary and

$$
\eta=\eta(\varepsilon):=\frac{\varepsilon \mathrm{e}^{-\int_{h}^{+\infty} \varphi(s) \mathrm{d} s} \mathrm{e}^{-D h}}{\lambda^{2} D_{1} \sqrt{\beta(0)+\delta(0)+2} \sqrt{\beta(h)+\delta(h)+2}} .
$$

Then, if $\left\|z_{0}\right\|_{0}<\eta$, by (2.19) and the hypothesis (H4) it results

$$
\begin{equation*}
\|z(t)\|_{0} \leq \frac{\varepsilon \mathrm{e}^{-\int_{h}^{+\infty} \varphi(s) \mathrm{d} s}}{\lambda \sqrt{\beta(h)+\delta(h)+2}}, \quad \forall t \in\left[t_{0}, h\right] . \tag{2.24}
\end{equation*}
$$

From the relations (2.23), (2.24), and the hypothesis (H4), it follows that $\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0}<\varepsilon$, $\forall t \geq h$.

Case 2: $t_{0} \geq h$. We similarly get

$$
\begin{equation*}
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0} \leq \lambda \sqrt{\beta\left(t_{0}\right)+\delta\left(t_{0}\right)+2}\left\|z_{0}\right\|_{0} \mathrm{e}^{-K \int_{t_{0}}^{t} \tilde{f}(s) \mathrm{d} s} \mathrm{e}^{\mathrm{e}_{t_{0}}^{t} \varphi(s) \mathrm{d} s} \tag{2.25}
\end{equation*}
$$

for all $t \geq t_{0}$. With the same $\eta$ as before, if $\left\|z_{0}\right\|_{0}<\eta$, then $\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0}<\varepsilon, \forall t \geq t_{0}$.
Therefore, the null solution of (1.1) is uniformly stable.
Proof of b). If, in addition (H3) holds, then from (2.25) we can easily obtain that the null solution of (1.1) is asymptotically stable.
Proof of $c$ ). We know from a) that the null solution of (1.1) is uniformly stable. It remains to prove that there exists $\xi>0$, such that for every $\varepsilon>0$ there exists $T=T(\varepsilon)>0$, such that $\left\|z_{0}\right\|_{0}<\xi$ implies $\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0}<\varepsilon$, for all $t_{0} \geq 0$ and $t \geq t_{0}+T$.

Indeed, if (H7) also holds, then $\int_{t_{0}}^{t} \widetilde{f}(s) \mathrm{d} s \geq p\left(t-t_{0}\right), \forall t \geq t_{0} \geq 0$. From (2.25) we obtain for all $t \geq t_{0} \geq 0$, that

$$
\begin{equation*}
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0} \leq \lambda \sqrt{\beta(0)+\delta(0)+2}\left\|z_{0}\right\|_{0} \mathrm{e}^{-K p\left(t-t_{0}\right)} N \tag{2.26}
\end{equation*}
$$

where $N:=\mathrm{e}^{\int_{0}^{+\infty} \varphi(s) \mathrm{d} s}$. Let $\xi:=\frac{1}{\lambda \sqrt{\beta(0)+\delta(0)+2}}, \varepsilon>0$, and

$$
T=T(\varepsilon):=\left\{\begin{array}{l}
\frac{1}{K p} \ln \frac{N}{\varepsilon}, \text { if } \varepsilon<N, \\
0, \text { if } \varepsilon \geq N .
\end{array}\right.
$$

Consider $z_{0} \in \mathbb{R}^{4}, z_{0} \neq 0$, with $\left\|z_{0}\right\|_{0}<\xi$ and let $t_{0} \geq 0$. Then for all $t \geq t_{0}+T$, by (2.26) we successively deduce

$$
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0}<\lambda \sqrt{\beta(0)+\delta(0)+2} \xi \mathrm{e}^{-K p\left(t-t_{0}\right)} N=N \mathrm{e}^{-K p\left(t-t_{0}\right)} \leq \varepsilon .
$$

Therefore the null solution of (1.1) is uniformly asymptotically stable.
Example 2.4. An example of functions $f_{i}, f_{j}, \beta, \delta, \gamma_{i}, g_{i}, i \in\{1,2\}, j \in\{3,4\}$, is

$$
\begin{gathered}
f_{1}(t)=\frac{1}{2 t+\sqrt{t^{2}+2}}, \quad f_{2}(t)=\frac{1}{t+\sqrt{t^{2}+1}}, \quad f_{3}(t)=\frac{1}{(t+1)^{4}}, \quad f_{4}(t)=\frac{2}{(t+1)^{3}}, \quad \forall t \geq 0, \\
\beta(t)=\frac{2 t+3}{t+1}, \quad \delta(t)=\frac{2 t^{3}+5}{t^{3}+2}, \quad \gamma_{1}(t)=\frac{1}{t \sqrt{t^{2}+1}+1}, \quad \gamma_{2}(t)=\mathrm{e}^{-t / 2}, \quad \forall t \geq 0, \\
g_{1}(t, x, y)=\mathrm{e}^{-t^{2} / 2} x^{3}, \quad g_{2}(t, x, y)=\frac{3}{t^{2} \sqrt{t}+1} y^{4}, \quad \forall t \geq 0, \quad \forall x, y \in \mathbb{R} .
\end{gathered}
$$

These functions satisfy the hypotheses (H1)-(H6), with $\beta_{0}=2, \delta_{0}=2, K_{1}=1 / \sqrt{2}, K_{2}=$ $(2+\sqrt{3}) \times(3-2 \sqrt{2}), h=1, r_{1}(t)=\mathrm{e}^{-t^{2} / 2}, r_{2}(t)=\frac{3}{t^{2} \sqrt{t+1}}, \forall t \geq 0$. In Figure 2.1 the solution of (1.1) and its derivative are plotted on two time intervals, for small initial data. The solution in the planes $(x, \dot{x})$ and $(y, \dot{y})$ on the same time intervals can be observed in Figure 2.2.

Example 2.5. If in Example 2.4 one changes only $f_{1}, f_{2}$ to $f_{1}(t)=\frac{1}{10}+\frac{1}{t+1}$, respectively $f_{2}(t)=\frac{1}{5}+\frac{2}{t+1}, \forall t \geq 0$, then the hypotheses $(\mathrm{H} 1),(\mathrm{H} 2),(\mathrm{H} 4)-(\mathrm{H} 7)$ are verified with $K_{1}=1 / 5$, $K_{2}=4 / 5, h=7, p=\frac{1}{10}$, and the same $\beta_{0}, \delta_{0}, r_{1}(t), r_{2}(t)$ and we obtain the solution of (1.1) and its derivative plotted in Figure 2.3 on the same time intervals and for the same initial data. In Figure 2.4 the solution is generated in the planes $(x, \dot{x})$ and $(y, \dot{y})$.


Figure 2.1: The solution of system (1.1) and its derivative, with the initial data $z_{0}=[0.01,0.01,0.01,0.01]$ and the functions $f_{1}, f_{2}, f_{3}, f_{4}, \beta, \delta, \gamma_{1}, \gamma_{2}, g_{1}, g_{2}$ given in Example 2.4.


Figure 2.2: The solution of (1.1) in the planes $(x, \dot{x})$ and $(y, \dot{y})$, with the data from Example 2.4.


Figure 2.3: The solution of system (1.1) and its derivative, with the initial data $z_{0}=[0.01,0.01,0.01,0.01]$ and the functions $f_{1}, f_{2}, f_{3}, f_{4}, \beta, \delta, \gamma_{1}, \gamma_{2}, g_{1}, g_{2}$ given in Example 2.5.


Figure 2.4: The solution of (1.1) in the planes $(x, \dot{x})$ and $(y, \dot{y})$, with the data from Example 2.5.

### 2.2. A stability result via Lyapunov's method

We are going to use the following additional assumptions.
$\left(\mathrm{H} 1^{*}\right) f_{i} \in C\left(\mathbb{R}_{+}\right) \cap L^{\infty}\left(\mathbb{R}_{+}\right), f_{j} \in C\left(\mathbb{R}_{+}\right), f_{i}(t) \geq 0, f_{j}(t) \geq 0, \forall t \in \mathbb{R}_{+}$, and $\int_{0}^{+\infty} f_{j}(t) \mathrm{d} t<$ $+\infty, \forall i \in\{1,2\}, \forall j \in\{3,4\} ;$
$\left(\mathrm{H} 3^{*}\right) \int_{0}^{+\infty} \widetilde{f}(t) \mathrm{d} t<+\infty$;
$\left(\mathrm{H} 4^{*}\right) \beta, \delta \in \mathrm{C}^{1}\left(\mathbb{R}_{+}\right), \beta, \delta$ are decreasing and

$$
\beta(t) \geq \beta_{0}>0, \delta(t) \geq \delta_{0}>0, \quad \forall t \in \mathbb{R}_{+} .
$$

Let us state and prove the following result.
Theorem 2.6. Suppose that the hypotheses (H1*), (H3*), (H4*), (H5), (H6) are fulfilled. Then the null solution of the system (1.1) is uniformly stable.

Proof. Let us remark that using the classical change of variables $x=x, u=\dot{x}, y=y, v=\dot{y}$, the system (1.1) becomes

$$
\begin{equation*}
\dot{z}=F(t, z), \tag{2.27}
\end{equation*}
$$

where

$$
z=\left(\begin{array}{c}
x \\
u \\
y \\
v
\end{array}\right), \quad F(t, z)=\left(\begin{array}{c}
u \\
-\beta(t) x-2 f_{1}(t) u+\gamma_{1}(t) y+f_{3}(t) v-g_{1}(t, x, y) \\
v \\
\gamma_{2}(t) x+f_{4}(t) u-\delta(t) y-2 f_{2}(t) v-g_{2}(t, x, y)
\end{array}\right)
$$

and our stability question reduces to the stability of the null solution $z(t)=0$ of the system (2.27). Let us remark that the global existence in the future of the solutions of (2.27) follows as in the proof of Theorem 2.3, this time the boundedness of the functions $f_{1}, f_{2}$ being ensured by the hypothesis ( $\mathrm{H} 1^{*}$ ).

We are going to use again the norm $\|\cdot\|_{0}$ defined by (2.14). Consider the function $V$ : $\mathbb{R}_{+} \times \Delta \rightarrow \mathbb{R}$,

$$
V(t, z)=\frac{1}{2}\left[\beta(t) x^{2}+u^{2}+\delta(t) y^{2}+v^{2}\right] \mathrm{e}^{-\int_{0}^{t}\left[\tilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+(s)}{\sqrt{\zeta(s)}}\right] \mathrm{d} s},
$$

for $z=(x, u, y, v)^{\top} \in \Delta$, where $\Delta \subset \mathbb{R}^{4}$ is a neighborhood of the origin of $\mathbb{R}^{4}$,

$$
\Delta=\left\{z \in \mathbb{R}^{4},\|z\|_{0}<a\right\}
$$

where $a=\min \left\{a_{1} \sqrt{\beta_{0}}, a_{2} \sqrt{\delta_{0}}\right\}, a_{1}>0, a_{2}>0$ are as in Remark 2.2, $\gamma(t):=\max \left\{\gamma_{1}(t)\right.$, $\left.\gamma_{2}(t)\right\}, \zeta(t):=\min \{\beta(t), \delta(t)\}, \forall t \in \mathbb{R}_{+}$, and $r(t):=\max \left\{r_{1}(t), r_{2}(t)\right\}, \forall t \geq 0$.

Obviously,

$$
\begin{aligned}
V(t, z) & \geq \frac{1}{2}\left(\beta_{0} x^{2}+u^{2}+\delta_{0} y^{2}+v^{2}\right) \mathrm{e}^{-\int_{0}^{t}\left[\tilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+(s)}{\sqrt{\xi(s)}}\right] \mathrm{d} s} \\
& =\frac{1}{2}\|z\|_{0}^{2} \mathrm{e}^{-\int_{0}^{t}\left[\tilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+(s)}{\sqrt{\xi(s)}}\right] \mathrm{d} s},
\end{aligned}
$$

for all $(t, z) \in \mathbb{R}_{+} \times \Delta$.
By using hypotheses (H1*), (H3*), (H4*), (H5), (H6), we deduce

$$
V(t, z) \geq \frac{1}{2}\|z\|_{0}^{2} \mathrm{e}^{-\left[\int_{0}^{+\infty} \tilde{f}(s) \mathrm{d} s+\int_{0}^{+\infty} f_{3}(s) \mathrm{d} s+\int_{0}^{+\infty} f_{4}(s) \mathrm{d} s+\int_{0}^{+\infty} \frac{\gamma(s)+(s)}{\sqrt{\zeta}(s)} \mathrm{d} s\right]}, \quad \forall(t, z) \in \mathbb{R}_{+} \times \Delta
$$

and so the function $V$ is positive definite.
The function $V$ is also decrescent. Indeed,

$$
\begin{aligned}
V(t, z) & \leq \frac{1}{2}\left[\beta(0) x^{2}+u^{2}+\delta(0) y^{2}+v^{2}\right] \mathrm{e}^{-\int_{0}^{t}\left[\tilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+\gamma s}{\sqrt{\xi(s)}}\right] \mathrm{d} s} \\
& \leq \frac{1}{2} \max \left\{\frac{\beta(0)}{\beta_{0}}, \frac{\delta(0)}{\delta_{0}}\right\}\|z\|_{0}^{2}, \quad \forall(t, z) \in \mathbb{R}_{+} \times \Delta .
\end{aligned}
$$

We prove that the time derivative of $V$ along the solutions of the system (2.27) is less than
or equal to 0 . Indeed, for every $(t, z) \in \mathbb{R}_{+} \times \Delta$,

$$
\begin{align*}
\frac{\mathrm{d} V}{\mathrm{~d} t}(t, z)= & \frac{1}{2}\left[\dot{\beta}(t) x^{2}+2 \beta(t) x \dot{x}+2 u \dot{u}+\dot{\delta}(t) y^{2}+2 \delta(t) y \dot{y}+2 v \dot{v}\right] \\
& \times \mathrm{e}^{-\int_{0}^{t}\left[\tilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+(s)}{\sqrt{\zeta(s)}}\right] \mathrm{d} s} \\
& -\left[\widetilde{f}(t)+f_{3}(t)+f_{4}(t)+\frac{\gamma(t)+r(t)}{\sqrt{\zeta(t)}}\right] V(t, z) \\
\leq & \left\{\gamma(t)(|u||y|+|x||v|)+\left[f_{3}(t)+f_{4}(t)\right]|u||v|+|u|\left|g_{1}(t, x, y)\right|+|v|\left|g_{2}(t, x, y)\right|\right\} \\
& \times \mathrm{e}^{-\int_{0}^{t}\left[\tilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+(s)}{\sqrt{\zeta(s)}}\right] \mathrm{d} s} \\
& -2\left[f_{1}(t) u^{2}+f_{2}(t) v^{2}\right] \mathrm{e}^{-\int_{0}^{t}\left[\tilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+r(s)}{\sqrt{\zeta(s)}}\right] \mathrm{d} s} \\
& -\left[\widetilde{f}(t)+f_{3}(t)+f_{4}(t)+\frac{\gamma(t)+r(t)}{\sqrt{\zeta(t)}}\right] V(t, z) . \tag{2.28}
\end{align*}
$$

From (2.28) and (2.3) for all $(t, z) \in \mathbb{R}_{+} \times \Delta$ we successively obtain

$$
\begin{align*}
\frac{\mathrm{d} V}{\mathrm{~d} t}(t, z) \leq & \left\{\gamma(t)(|u||y|+|x||v|)+\left[f_{3}(t)+f_{4}(t)\right]|u||v|+\left[r_{1}(t)|x||u|+r_{2}(t)|y||v|\right]\right. \\
& \left.-2\left[f_{1}(t) u^{2}+f_{2}(t) v^{2}\right]\right\} \mathrm{e}^{-\int_{0}^{t}\left[\tilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+(s)}{\sqrt{\zeta}(s)}\right] \mathrm{d} s} \\
& -\left[\widetilde{f}(t)+f_{3}(t)+f_{4}(t)+2 \frac{\gamma(t)+r(t)}{\sqrt{\zeta(t)}}\right] V(t, z) \\
\leq & {\left[f_{3}(t)+f_{4}(t)+\frac{\gamma(t)+r(t)}{\sqrt{\zeta(t)}}\right] V-2\left[f_{1}(t) u^{2}+f_{2}(t) v^{2}\right] } \\
& \times \mathrm{e}^{-\int_{0}^{t}\left[\tilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+r(s)}{\sqrt{\zeta(s)}}\right] \mathrm{d} s} \\
& -\left[\widetilde{f}(t)+f_{3}(t)+f_{4}(t)+\frac{\gamma(t)+r(t)}{\sqrt{\zeta(t)}}\right] V(t, z) \\
= & -\widetilde{f}(t) V-2\left[f_{1}(t) u^{2}+f_{2}(t) v^{2}\right] \mathrm{e}^{-\int_{0}^{t}\left[\tilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+(s)}{\sqrt{\zeta(s)}}\right] \mathrm{d} s} . \tag{2.29}
\end{align*}
$$

Then, from (2.29) we easily get

$$
\frac{\mathrm{d} V}{\mathrm{~d} t}(t, z) \leq 0, \quad \forall(t, z) \in \mathbb{R}_{+} \times \Delta
$$

From Persidski's Theorem (see, e.g., [3, second Corollary, p. 101], [17, Theorem 2.1]), it follows that the null solution of (1.1) is uniformly stable.

Remark 2.7. Let us remark that by using the transformation (2.4) we obtained the uniform, the asymptotic, and the uniform asymptotic stability, while by using the classical transformation ( $x=x, u=\dot{x}, y=y, v=\dot{y}$ ) and the Lyapunov's method we were only able to achieve the uniform stability of the null solution of (1.1). Hence the first method, based on the transformation (2.4), is more effective.

Remark 2.8. Note that the null solution of the system (1.1) can be uniformly stable and not asymptotically stable. Indeed, this can be seen by considering the following functions

$$
\begin{gathered}
f_{1}(t)=\frac{\mathrm{e}^{-t}}{t+1}, \quad f_{2}(t)=\frac{\left|\cos ^{3} t\right|}{t^{2}+4}, \quad \forall t \geq 0, \quad f_{3}(t)=\frac{\left|\sin t^{2}\right|}{t+2}, \quad f_{4}(t)=\frac{\mathrm{e}^{-t^{2}}}{t+1}, \quad \forall t \geq 0 \\
\beta(t)=0.3+\frac{1}{t^{2}+1}, \quad \delta(t)=0.2+\frac{1}{\sqrt{t^{2}+2}}, \quad \gamma_{1}(t)=\frac{t}{t+2} \mathrm{e}^{-t^{2}}, \quad \gamma_{2}(t)=\frac{3|\cos t|}{(t+1)^{2}}, \quad \forall t \geq 0 \\
g_{1}(t, x, y)=\frac{3 x^{3}}{\left(t^{2}+2\right)^{2}}, \quad g_{2}(t, x, y)=\frac{2 y^{2}}{(t+1)^{3}}, \quad \forall t \geq 0, \quad \forall x, y \in \mathbb{R}
\end{gathered}
$$

These functions satisfy the hypotheses $\left(\mathrm{H} 1^{*}\right),\left(\mathrm{H} 3^{*}\right),\left(\mathrm{H} 4^{*}\right),(\mathrm{H} 5),(\mathrm{H} 6)$, with $\beta_{0}=0.3, \delta_{0}=0.2$, $r_{1}(t)=\frac{3}{\left(t^{2}+2\right)^{2}}, r_{2}(t)=\frac{2}{(t+1)^{3}}, \forall t \geq 0$. For small initial data, the solution of (1.1) and its derivative can be observed in Figure 2.5 on some time intervals. The plottings of the solution in the planes $(x, \dot{x}),(y, \dot{y})$ are given in Figure 2.6.


Figure 2.5: The solution of (1.1) and its derivative, with the initial data $z_{0}=$ $[0.001,0.001,0.001,0.001]$ and the functions $f_{1}, f_{2}, f_{3}, f_{4}, \beta, \delta, \gamma_{1}, \gamma_{2}, g_{1}, g_{2}$ given in Remark 2.8.


Figure 2.6: The solution of (1.1) in the planes $(x, \dot{x})$ and $(y, \dot{y})$, with the data from Remark 2.8.

## 3. Analysis of the inhomogeneous system (1.3)

Suppose that the block of mass $m_{1}$ is subject to the action of a time-dependent external force $\widehat{f}: \mathbb{R}_{+} \rightarrow \mathbb{R}$. In this case, we obtain the inhomogeneous system (1.3).

We are going to use the following hypotheses.
(H8) $f \in C\left(\mathbb{R}_{+}\right)$and $f \in L^{1}\left(\mathbb{R}_{+}\right)$;
(H9) $f \in C\left(\mathbb{R}_{+}\right)$and $\lim _{t \rightarrow+\infty} f(t)=0$.

### 3.1. Qualitative properties of solutions via differential inequalities

## Theorem 3.1.

a) Suppose that the hypotheses (H1), (H2), (H4)-(H6), (H8) are fulfilled. Then every solution of the system (1.3) starting from sufficiently small initial data and its derivative are bounded.
b) If the hypotheses (H1), (H2), (H4)-(H6), (H7) with p big enough, and (H9) are satisfied, then for every solution $(x, y)$ of (1.3) starting from small initial data, we have $\lim _{t \rightarrow+\infty} x(t)=$ $\lim _{t \rightarrow+\infty} \dot{x}(t)=\lim _{t \rightarrow+\infty} y(t)=\lim _{t \rightarrow+\infty} \dot{y}(t)=0$.

Proof. This time we use the following transformation (of the same type as the one from [2])

$$
\left\{\begin{array}{l}
\dot{x}=u-f_{1}(t) x  \tag{3.1}\\
\dot{u}=\left[\dot{f}_{1}(t)+f_{1}^{2}(t)-\beta(t)\right] x-f_{1}(t) u+\left[\gamma_{1}(t)-f_{2}(t) f_{3}(t)\right] y+f_{3}(t) v+f(t)-g_{1}(t, x, y) \\
\dot{y}=v-f_{2}(t) y \\
\dot{v}=\left[\gamma_{2}(t)-f_{1}(t) f_{4}(t)\right] x+f_{4}(t) u+\left[\dot{f}_{2}(t)+f_{2}^{2}(t)-\delta(t)\right] y-f_{2}(t) v-g_{2}(t, x, y)
\end{array}\right.
$$

and the system (1.3) becomes

$$
\begin{equation*}
\dot{z}=A(t) z+B(t) z+G(t, z), \tag{3.2}
\end{equation*}
$$

where

$$
G(t, z)=\left(\begin{array}{c}
0 \\
f(t)-g_{1}(t, x, y) \\
0 \\
-g_{2}(t, x, y)
\end{array}\right)
$$

and $A(t)$ and $B(t)$ are the same as in the proof of Theorem 2.3.
Let $z_{0} \in \mathbb{R}^{4} \backslash\{0\}$ with $\left\|z_{0}\right\|_{0}$ small enough, $t_{0} \geq 0$, and

$$
z\left(t, t_{0}, z_{0}\right)=\left(x\left(t, t_{0}, z_{0}\right), u\left(t, t_{0}, z_{0}\right), y\left(t, t_{0}, z_{0}\right), v\left(t, t_{0}, z_{0}\right)\right)^{\top}
$$

be the unique solution of (3.2) which is equal to $z_{0}$ for $t=t_{0}$.
Similarly (by applying, e.g., [3, Corollary, p. 53]) we conclude that $z\left(t, t_{0}, z_{0}\right)$ exists on $\left[t_{0},+\infty\right)$, this time having

$$
\|A(t) z+B(t) z+G(t, z)\|_{0} \leq \psi(t)\|z\|_{0}+|f(t)|, \quad \forall(t, z) \in \mathbb{R}_{+} \times \mathbb{R}^{4} .
$$

As before we deduce

$$
\begin{align*}
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0} \leq & \left.\lambda\left\|z_{0}\right\|_{0} \sqrt{\beta\left(t_{0}\right)+\delta\left(t_{0}\right)+2} \mathrm{e}^{\int_{t_{0}}^{t}\left[-\tilde{f}(u)+\frac{\gamma(u)}{2 \sqrt{\xi(u)}}\right]}\right] \mathrm{d} u \\
& \times \int_{t_{0}}^{t} \mathrm{e}^{\int_{s}^{t}\left[-\tilde{f}(u)+\frac{\gamma(u)}{2 \sqrt{\xi(u)}}\right] \mathrm{d} u} \\
& +f_{1}(s) f_{4}(s)\left|x\left(s, t_{0}, z_{0}\right)\right|+f_{2}^{2}(s) f_{3}(s)\left|y\left(s, t_{0}, z_{0}\right)\right| \\
& +f_{3}(s)\left|v\left(s, t_{0}, z_{0}\right)\right|+f_{4}(s)\left|u\left(s, t_{0}, z_{0}\right)\right|+|f(s)| \\
& +\left|g_{1}\left(s, x\left(s, t_{0}, z_{0}\right), y\left(s, t_{0}, z_{0}\right)\right)\right| \\
& \left.+\left|g_{2}\left(s, x\left(s, t_{0}, z_{0}\right), y\left(s, t_{0}, z_{0}\right)\right)\right|\right] \mathrm{d} s, \tag{3.3}
\end{align*}
$$

for all $t \geq t_{0}$.
We distinguish two cases again.
Case 1: $0 \leq t_{0}<h$. As in the proof of Theorem 2.3, we obtain the relation (2.19), with $D$, $D_{1}>0$.

From (3.3) and using Remark 2.2, we deduce for all $t \geq h$

$$
\begin{aligned}
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0} \leq & \lambda \sqrt{\beta(h)+\delta(h)+2}\left\|z\left(h, t_{0}, z_{0}\right)\right\|_{0} \mathrm{e}^{f_{h}^{t}}\left[-\tilde{f}(u)+\frac{\gamma(u)}{2 \sqrt{\zeta(u)}}\right] \mathrm{d} u \\
& +\int_{h}^{t} \mathrm{e}^{\int_{s}^{t}\left[-\tilde{f}(u)+\frac{\gamma(u)}{2 \sqrt{\zeta(u)}}\right] \mathrm{d} u\left\{\left[\left(\frac{K_{1}}{\sqrt{\beta_{0}}}+\frac{K_{2}}{\sqrt{\delta_{0}}}\right) \widetilde{f}(s)\right.\right.} \\
& +\frac{f_{1}(s) f_{4}(s)}{\sqrt{\beta_{0}}}+\frac{f_{2}(s) f_{3}(s)}{\sqrt{\delta_{0}}}+f_{3}(s)+f_{4}(s) \\
& \left.\left.+\frac{r_{1}(s)}{\sqrt{\beta_{0}}}+\frac{r_{2}(s)}{\sqrt{\delta_{0}}}\right]\left\|z\left(s, t_{0}, z_{0}\right)\right\|_{0}+|f(s)|\right\} \mathrm{d} s \\
=: & \rho(t), \quad \forall t \geq h .
\end{aligned}
$$

Straightforward calculations lead us to

$$
\left\{\begin{array}{l}
\dot{\rho}(t) \leq \omega(t) \rho(t)+|f(t)|, \quad \forall t \geq h \\
\rho(h)=\lambda \sqrt{\beta(h)+\delta(h)+2}\left\|z\left(h, t_{0}, z_{0}\right)\right\|_{0}
\end{array}\right.
$$

with $\omega(t), t \geq 0$, as in the proof of Theorem 2.3.
We easily deduce

$$
\begin{align*}
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0} & \leq\left(\rho(h)+\int_{h}^{t} \mathrm{e}^{-\int_{h}^{s}[-K \tilde{f}(u)+\varphi(u)] \mathrm{d} u}|f(s)| \mathrm{d} s\right) \mathrm{e}^{\int_{h}^{t}[-\kappa \tilde{f}(s)+\varphi(s)] \mathrm{d} s} \\
& =: \mu(t), \quad \forall t \geq h . \tag{3.4}
\end{align*}
$$

Proof of a). By using the hypotheses (H1), (H5), (H6), and Remark 2.1, it is readily seen that $\bar{\varphi}:=\int_{0}^{+\infty} \varphi(t) \mathrm{d} t<+\infty$. From (3.4) and the hypothesis (H8) we derive that

$$
\begin{aligned}
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0} & \leq \rho(h) \mathrm{e}^{\int_{h}^{t} \varphi(u) \mathrm{d} u}+\int_{h}^{t} \mathrm{e}^{\int_{s}^{t} \varphi(u) \mathrm{d} u}|f(s)| \mathrm{d} s \\
& \leq \mathrm{e}^{\overline{\bar{\varphi}}}\left(\rho(h)+\int_{h}^{t}|f(s)| \mathrm{d} s\right) \\
& \leq \mathrm{e}^{\bar{\varphi}}\left(\rho(h)+\|f\|_{L^{1}[0,+\infty)}\right)<+\infty, \quad \forall t \geq h
\end{aligned}
$$

and so every solution of (1.3) with initial data small enough is bounded. The boundedness of $\dot{z}\left(t, t_{0}, z_{0}\right)$ follows immediately.

Proof of $b$ ). Let us estimate the limit of $\mu$ at $+\infty$. We have

$$
\begin{equation*}
\lim _{t \rightarrow+\infty} \mu(t)=\lim _{t \rightarrow+\infty} \frac{\rho(h)+\int_{h}^{t} \mathrm{e}^{-\int_{h}^{s}[-K \tilde{f}(u)+\varphi(u)] \mathrm{d} u}|f(s)| \mathrm{d} s}{\mathrm{e}^{-\int_{h}^{t}[-K \widetilde{f}(s)+\varphi(s)] \mathrm{d} s}} \tag{3.5}
\end{equation*}
$$

If $\int_{h}^{+\infty} \mathrm{e}^{-\int_{h}^{s}[-K \tilde{f}(u)+\varphi(u)] \mathrm{d} u}|f(s)| \mathrm{d} s<+\infty$, then, from (3.5) and the hypothesis (H7), we easily obtain

$$
\lim _{t \rightarrow+\infty} \mu(t)=0
$$

If $\int_{h}^{+\infty} \mathrm{e}^{-\int_{h}^{s}[-K \tilde{f}(u)+\varphi(u)] \mathrm{d} u}|f(s)| \mathrm{d} s=+\infty$, then we estimate

$$
\begin{equation*}
\lim _{t \rightarrow+\infty} \frac{\frac{\mathrm{d}}{\mathrm{~d} t}\left(\rho(h)+\int_{h}^{t} \mathrm{e}^{-\int_{h}^{s}[-K \widetilde{f}(u)+\varphi(u)] \mathrm{d} u}|f(s)| \mathrm{d} s\right)}{\frac{\mathrm{d}}{\mathrm{~d} t}\left(\mathrm{e}^{-\int_{h}^{t}[-K \widetilde{f}(s)+\varphi(s)] \mathrm{d} s}\right)}=\lim _{t \rightarrow+\infty} \frac{|f(t)|}{K \widetilde{f}(t)-\varphi(t)} \tag{3.6}
\end{equation*}
$$

Using the hypotheses (H1), (H5)-(H7), and Remark 2.1,

$$
K \widetilde{f}(t)-\varphi(t) \geq K p-\varphi_{0}, \quad \forall t \geq 0
$$

where $\varphi_{0}=\sup _{t \geq 0}\{\varphi(t)\}$. Hence, if $p>\frac{\varphi_{0}}{K}$, then $K \widetilde{f}(t)-\varphi(t)>0, \forall t \geq 0$, and, from (3.6), the hypothesis (H9), and L'Hospital's rule, we obtain $\lim _{t \rightarrow+\infty} \mu(t)=0$. Hence, by (3.4) it follows that $\lim _{t \rightarrow+\infty}\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0}=0$ and we also infer $\lim _{t \rightarrow+\infty}\left\|\dot{z}\left(t, t_{0}, z_{0}\right)\right\|_{0}=0$.
Case 2: $t_{0} \geq h$. The proofs of a ) and b ) follow as in Case 1 , this time by using the inequality

$$
\begin{aligned}
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0} \leq & \left(\lambda \sqrt{\beta\left(t_{0}\right)+\delta\left(t_{0}\right)+2}\left\|z_{0}\right\|_{0}+\int_{t_{0}}^{t} \mathrm{e}^{-\int_{t_{0}}^{s}[-K \tilde{f}(u)+\varphi(u)] \mathrm{d} u}|f(s)| \mathrm{d} s\right) \\
& \times \mathrm{e}^{\int_{t_{0}}^{t}[-K \tilde{f}(s)+\varphi(s)] \mathrm{d} s}, \quad \forall t \geq t_{0}
\end{aligned}
$$

Example 3.2. If we consider the functions

$$
\begin{gathered}
f_{1}(t)=\left\{\begin{array}{l}
\frac{\ln t}{t}, t \geq \mathrm{e} \\
\frac{t}{\mathrm{e}^{3}}(2 \mathrm{e}-t), \quad t \in[0, \mathrm{e})^{2}
\end{array}, \quad f_{2}(t)=\left\{\begin{array}{l}
\frac{\ln t}{t-1}, t \geq \mathrm{e} \\
\frac{t}{\mathrm{e}(\mathrm{e}-1)^{2}}(2 \mathrm{e}-1-t), \quad t \in[0, \mathrm{e})^{\prime}
\end{array}\right.\right. \\
f_{3}(t)=\frac{\arctan t}{(t+1)^{2}}, \quad f_{4}(t)=\frac{\sqrt{t}}{(t+2)^{2}}, \quad f(t)=\frac{2 t+3}{t+2} \mathrm{e}^{-t}, \quad \forall t \geq 0, \\
\beta(t)=\frac{9}{\mathrm{e}^{2}}+\frac{1}{\sqrt{t+2}}, \delta(t)=\frac{49}{4(\mathrm{e}-1)^{2}}+\mathrm{e}^{-2 t}, \gamma_{1}(t)=\frac{\mathrm{e}^{-3 t}}{t^{2}+1}, \quad \gamma_{2}(t)=\frac{\sin ^{2} t}{(t+1)^{3}}, \quad \forall t \geq 0, \\
g_{1}(t, x, y)=\frac{2|\sin t| x^{3}}{t \sqrt{t}+1}, \quad g_{2}(t, x, y)=\frac{3|\cos t| y^{2}}{(t+1) \sqrt{t+1}}, \quad \forall t \geq 0, \quad \forall x, y \in \mathbb{R},
\end{gathered}
$$

then the hypotheses (H1), (H2), (H4)-(H6), (H8) are fulfilled with $\beta_{0}=\frac{9}{\mathrm{e}^{2}}, \delta_{0}=\frac{49}{4(\mathrm{e}-1)^{2}}$, $K_{1}=2 / \mathrm{e}, K_{2}=1 /(\mathrm{e}-1), h=\mathrm{e}, r_{1}(t)=\frac{2|\sin t|}{t \sqrt{t}+1}, r_{2}(t)=\frac{3|\cos t|}{(t+1) \sqrt{t+1}}, \forall t \geq 0$. In Figure 3.1 one can observe the solution of (1.3) and its derivative, for small initial data on two time intervals and in Figure 3.2 the solution is plotted in the planes $(x, \dot{x}),(y, \dot{y})$ on the same time intervals.


Figure 3.1: The solution of (1.3) and its derivative, with the initial data $z_{0}=$ $[0.01,0.01,0.01,0.01]$ and the functions $f_{1}, f_{2}, f_{3}, f_{4}, f, \beta, \delta, \gamma_{1}, \gamma_{2}, g_{1}, g_{2}$ given in Example 3.2.


Figure 3.2: The solution of (1.3) in the planes $(x, \dot{x})$ and $(y, \dot{y})$, with the data from Example 3.2.

Remark 3.3. Let us remark the difference between the graphs of the first and second components of the solution near the origin. Due to the action of the external force $\widehat{f}(t)$ on the first block $m_{1}$, at least near the origin, the absolute values of $x=x(t)$ are much bigger than the ones of $y=y(t)$.

### 3.2. Boundedness of solutions

Theorem 3.4. Suppose that the hypotheses (H1*), (H4*), (H5), (H6), (H8) are fulfilled. Then every solution of the system (1.3) with sufficiently small initial data is bounded.

Proof. Let us remark that using the classical change of variables $x=x, u=\dot{x}, y=y, v=\dot{y}$, the system (1.3) becomes

$$
\begin{equation*}
\dot{z}=F(t, z), \tag{3.7}
\end{equation*}
$$

where

$$
z=\left(\begin{array}{l}
x \\
u \\
y \\
v
\end{array}\right) \quad F(t, z)=\left(\begin{array}{c}
u \\
-\beta(t) x-2 f_{1}(t) u+\gamma_{1}(t) y+f_{3}(t) v+f(t)-g_{1}(t, x, y) \\
v \\
\gamma_{2}(t) x+f_{4}(t) u-\delta(t) y-2 f_{2}(t) v-g_{2}(t, x, y)
\end{array}\right) .
$$

We will use again the norm $\|\cdot\|_{0}$ defined by (2.14) and the function $V: \mathbb{R}_{+} \times \Delta \rightarrow \mathbb{R}$,

$$
V(t, z)=\frac{1}{2}\left[\beta(t) x^{2}+u^{2}+\delta(t) y^{2}+v^{2}\right] \mathrm{e}^{-\int_{0}^{t}\left[\tilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+r(s)}{\sqrt{\xi(s)}}\right] \mathrm{d} s}
$$

for $z=(x, u, y, v)^{\top} \in \Delta$, where $\Delta \subset \mathbb{R}^{4}$ is as in the proof of Theorem 2.6, $\gamma(t):=\max \left\{\gamma_{1}(t)\right.$, $\left.\gamma_{2}(t)\right\}, \zeta(t):=\min \{\beta(t), \delta(t)\}, \forall t \in \mathbb{R}_{+}$, and $r(t):=\max \left\{r_{1}(t), r_{2}(t)\right\}, \forall t \geq 0$.

Let us calculate the time derivative of $V$ along the solutions of the system (3.7), whose global existence in the future is deduced as in the proof of Theorem 2.6. For every $(t, z) \in$ $\mathbb{R} \times \Delta$ we have

$$
\begin{aligned}
\frac{\mathrm{d} V}{\mathrm{~d} t}(t, z)= & \frac{1}{2}\left[\dot{\beta}(t) x^{2}+2 \beta(t) x \dot{x}+2 u \dot{u}+\dot{\delta}(t) y^{2}+2 \delta(t) y \dot{y}+2 v \dot{v}\right] \\
& \times \mathrm{e}^{-\int_{0}^{t}\left[\tilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+(s)}{\sqrt{\zeta(s)}}\right] \mathrm{d} s} \\
& -\left[\widetilde{f}(t)+f_{3}(t)+f_{4}(t)+\frac{\gamma(t)+r(t)}{\sqrt{\zeta(t)}}\right] V(t, z) .
\end{aligned}
$$

By hypothesis ( $\mathrm{H} 4^{*}$ ) we get for every $(t, z) \in \mathbb{R}_{+} \times \Delta$,

$$
\begin{align*}
\frac{\mathrm{d} V}{\mathrm{~d} t}(t, z) \leq & \left\{\gamma(t)(|u||y|+|x||v|)+\left[f_{3}(t)+f_{4}(t)\right]|u||v|+|u|\left|g_{1}(t, x, y)\right|+|v|\left|g_{2}(t, x, y)\right|\right. \\
& +|f(t)||u|\} \times \mathrm{e}^{-\int_{0}^{t}\left[\widetilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+r(s)}{\sqrt{\zeta(s)}}\right] \mathrm{d} s} \\
& -\left[\widetilde{f}(t)+f_{3}(t)+f_{4}(t)+\frac{\gamma(t)+r(t)}{\sqrt{\zeta(t)}}\right] V(t, z) \\
& -2\left[f_{1}(t) u^{2}+f_{2}(t) v^{2}\right] \mathrm{e}^{-\int_{0}^{t}\left[\tilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+(s)}{\sqrt{\zeta(s)}}\right] \mathrm{d} s} . \tag{3.8}
\end{align*}
$$

From relations (3.8) and Remark 2.2, we successively deduce

$$
\begin{align*}
\frac{\mathrm{d} V}{\mathrm{~d} t}(t, z) \leq & \left\{\gamma(t)(|u||y|+|x||v|)+\left[f_{3}(t)+f_{4}(t)\right]|u||v|+\left[r_{1}(t)|x||u|+r_{2}(t)|y||v|\right.\right. \\
& \left.+|f(t)||u|]-2\left[f_{1}(t) u^{2}+f_{2}(t) v^{2}\right]\right\} \mathrm{e}^{-\int_{0}^{t}\left[\widetilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+r(s)}{\sqrt{\zeta(s)}}\right] \mathrm{d} s} \\
& -\left[\widetilde{f}(t)+f_{3}(t)+f_{4}(t)+\frac{\gamma(t)+r(t)}{\sqrt{\zeta(t)}}\right] V(t, z) \\
\leq & {\left[f_{3}(t)+f_{4}(t)+\frac{\gamma(t)+r(t)}{\sqrt{\zeta(t)}}\right] V(t, z)+|f(t)||u| \mathrm{e}^{-\int_{0}^{t}\left[\widetilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+(s)}{\sqrt{\zeta(s)}}\right] \mathrm{d} s} } \\
& -2\left[f_{1}(t) u^{2}+f_{2}(t) v^{2}\right] \mathrm{e}^{-\int_{0}^{t}\left[\tilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+(s)}{\sqrt{\zeta(s)}}\right] \mathrm{d} s} \\
& -\left[\widetilde{f}(t)+f_{3}(t)+f_{4}(t)+\frac{\gamma(t)+r(t)}{\sqrt{\zeta(t)}}\right] V(t, z) \\
\leq & -\tilde{f}(t) V(t, z)+|f(t)||u| \mathrm{e}^{-\int_{0}^{t}\left[\tilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+r(s)}{\sqrt{\zeta(s)}}\right] \mathrm{d} s}, \tag{3.9}
\end{align*}
$$

for all $(t, z) \in \mathbb{R}_{+} \times \Delta$. Then, from (3.9) we easily obtain $\forall(t, z) \in \mathbb{R}_{+} \times \Delta$

$$
\begin{aligned}
\frac{\mathrm{d} V}{\mathrm{~d} t}(t, z) \leq & -\widetilde{f}(t) V(t, z)+|f(t)| \sqrt{\beta(t) x^{2}+u^{2}+\delta(t) y+v^{2}} \\
& \times \mathrm{e}^{-\int_{0}^{t}\left[\widetilde{f}(s)+f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+(s)}{\sqrt{\zeta}(s)}\right] \mathrm{d} s} \\
\leq & -\widetilde{f}(t) V(t, z)+|f(t)| \sqrt{2 V(t, z)} \mathrm{e}^{-\frac{1}{2} \int_{0}^{t} \tilde{f}(s) \mathrm{d} s},
\end{aligned}
$$

which actually represents an inequality of Bernoulli type.
Let $z_{0} \in \Delta, t_{0} \geq 0$, and $z\left(t, t_{0}, z_{0}\right)$ be the unique solution of (3.7) which is equal to $z_{0}$ for $t=t_{0}$. Using classical differential estimates, we find

$$
V\left(t, z\left(t, t_{0}, z_{0}\right)\right) \leq \mathrm{e}^{-\int_{t_{0}}^{t} \tilde{f}(s) \mathrm{d} s}\left[\sqrt{V\left(t_{0}, z_{0}\right)}+\frac{\sqrt{2}}{2} \int_{t_{0}}^{t}|f(s)| \mathrm{e}^{-\frac{1}{2} \int_{0}^{t_{0}} \tilde{f}(u) \mathrm{d} u} \mathrm{~d} s\right]^{2}, \quad \forall t \geq t_{0}
$$

Therefore, by using the hypotheses ( $\mathrm{H} 1^{*}$ ), (H5), (H6), it follows that

$$
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0} \leq M\left[\sqrt{V\left(t_{0}, z_{0}\right)}+\frac{\sqrt{2}}{2} \int_{t_{0}}^{t}|f(s)| \mathrm{e}^{-\frac{1}{2} \int_{0}^{t_{0}} \tilde{f}(u) \mathrm{d} u} \mathrm{~d} s\right], \quad \forall t \geq t_{0}
$$

where $M:=\sqrt{2} \mathrm{e}^{\frac{1}{2} \int_{0}^{t_{0}} \tilde{f}(s) \mathrm{d} s+\frac{1}{2} \int_{0}^{+\infty}\left[f_{3}(s)+f_{4}(s)+\frac{\gamma(s)+(s)}{\sqrt{\xi(s)}}\right] \text { ds }}$. If the hypothesis (H8) comes into play, then

$$
\left\|z\left(t, t_{0}, z_{0}\right)\right\|_{0} \leq M\left[\sqrt{V\left(t_{0}, z_{0}\right)}+\frac{\sqrt{2}}{2}\|f\|_{L^{1}[0,+\infty)} \mathrm{e}^{-\frac{1}{2} \int_{0}^{t_{0}} \tilde{f}(s) \mathrm{d} s}\right], \quad \forall t \geq t_{0} .
$$

Remark 3.5. Note that by using the classical transformation ( $x=x, u=\dot{x}, y=y, v=\dot{y}$ ), we could only deduce the boundedness of the solutions of (1.3) for initial data small enough. In contrast, the transformation (3.1) allowed us to obtain in addition that the solutions of (1.3), starting from sufficiently small initial data, have the limit zero at $+\infty$.

## Acknowledgements

We are grateful for the remarks and suggestions of the anonymous reviewer and editor Bo Zhang, which led to an improved version of the paper.

## References

[1] R. Bellman, Stability theory of differential equations, McGraw-Hill, New York, 1953. MR0061235; Zbl 0053.24705
[2] T. A. Burton, T. Furumochi, A note on stability by Schauder's theorem, Funkcial. Ekvac. 44(2001), No. 1, 73-82. MR1847837; Zbl 1158.34329
[3] C. Corduneanu, Principles of differential and integral equations, Allyn and Bacon, Boston, Mass., 1971. MR0276520; Zbl 0208.10701
[4] R. R. Craig Jr., A. J. Kurdila, Fundamentals of structural dynamics, Wiley, Hoboken, 2006. Zbl 1118.70001
[5] J. K. Hale, Ordinary differential equations [2nd edition], Krieger, Florida, 1980. MR0587488; Zbl 0433.34003
[6] P. Hartman, Ordinary differential equations, Wiley, New York, 1964. MR0171038; Zbl 0125.32102
[7] L. Hatvani, On the stability of the zero solution of certain second order non-linear differential equations, Acta Sci. Math. (Szeged) 32(1971), 1-9. MR0306639; Zbl 0216.11704
[8] L. Hatvani, On the asymptotic behaviour of the solutions of $\left(p(t) x^{\prime}\right)^{\prime}+q(t) f(x)=0$, Publ. Math. Debrecen 19(1972), 225-237 (1973). MR0326064; Zbl 0271.34061
[9] L. Hatvani, A generalization of the Barbashin-Krasovskij theorems to the partial stability in nonautonomous systems, Qualitative theory of differential equations, Vol. I, II (Szeged, 1979), pp. 381-409, Colloq. Math. Soc. János Bolyai, 30, North-Holland, Amsterdam-New York, 1981. MR0680604; Zbl 0486.34044
[10] L. Hatvani, T. Krisztin, V. Totik, A necessary and sufficient condition for the asymptotic stability of the damped oscillator, J. Differential Equations 119(1995), No. 1, 209-223. https : //doi.org/10.1006/jdeq.1995.1087; MR1334491; Zbl 0831.34052
[11] L. Hatvani, Integral conditions on the asymptotic stability for the damped linear oscillator with small damping, Proc. Amer. Math. Soc. 124(1996), No. 2, 415-422. https: //doi.org/10.1090/S0002-9939-96-03266-2; MR1317039; Zbl 0844.34051
[12] S. H. Ju, H. H. Kuo, S. W. Yu, Investigation of vibration induced by moving cranes in high-tech factories, J: Low Frequency Noise, Vibration and Active Control 39(2019), No. 1, 84-97. https://doi.org/10.1177/1461348419837416
[13] V. Lakshmikantham, S. Leela, Differential and integral inequalities. Theory and applications, Vol. I: Ordinary differential equations, Mathematics in Science and Engineering, Vol. 55-I. Academic Press, New York-London, 1969. MR0379933; Zbl 0177.12403
[14] G. Moroșanu, C. Vladimirescu, Stability for a nonlinear second order ODE, Funkcial. Ekvac. 48(2005), No. 1, 49-56. https://doi.org/10.1619/fesi.48.49; MR2154377; Zbl 1122.34324
[15] G. Moroșanu, C. Vladimirescu, Stability for a damped nonlinear oscillator, Nonlinear Anal. 60(2005), No. 2, 303-310. https://doi.org/10.1016/j.na.2004.08.027; MR2101880; Zbl 1071.34049
[16] G. Moroșanu, C. Vladimirescu, Stability for a system of two coupled nonlinear oscillators with partial lack of damping, Nonlinear Anal. Real World Appl. 45(2019), 609-619. https://doi.org/10.1016/j.nonrwa.2018.07.026; MR3854325; Zbl 1418.34085
[17] G. Moroșanu, C. Vladimirescu, Stability for systems of 1-D coupled nonlinear oscillators, Nonlinear Anal. Real World Appl. 59(2021) 103242. https://doi.org/10.1016/j. nonrwa.2020.103242; MR4170813; Zbl 1469.34055
[18] M. Paz, W. Leigh, Structural dynamics: theory and computation [5th edition], Springer, New York, 2004. https://doi.org/10.1007/978-1-4615-0481-8
[19] L. C. Piccinini, G. Stampacchia, G. Vidossich, Ordinary differential dquations in $\mathbb{R}^{n}$. Problems and methods, Applied Mathematical Sciences, Vol. 39, Springer-Verlag, New York, 1984. https://doi.org/10.1007/978-1-4612-5188-0; MR0740539; Zbl 0535.34001
[20] P. Pucci, J. Serrin, Precise damping conditions for global asymptotic stability for nonlinear second order systems, Acta Math. 170(1993), No. 2, 275-307. https://doi.org/10. 1007/BF02392788; MR1226530; Zbl 0797.34059
[21] P. Pucci, J. Serrin, Precise damping conditions for global asymptotic stability for nonlinear second order systems, II, J. Differential Equations 113(1994), No. 2, 505-534. https : //doi.org/10.1006/jdeq.1994.1134; MR1297668; Zbl 0814.34033
[22] P. Pucci, J. Serrin, Asymptotic stability for intermittently controlled nonlinear oscillators, SIAM J. Math. Anal. 25(1994), No. 3, 815-835. https://doi.org/10.1137/ S0036141092240679; MR1271312; Zbl 0809.34067
[23] P. Pucci, J. Serrin, Asymptotic stability for ordinary differential systems with time dependent restoring potentials, Arch. Rational Mech. Anal. 132(1995), No. 3, 207-232. https://doi.org/10.1007/BF00382747; MR1365829; Zbl 0861.34034
[24] R. A. Smith, Asymptotic stability of $x^{\prime \prime}+a(t) x^{\prime}+x=0$, Quart. J. Math. Oxford Ser. 12(1961), No. 2, 123-126. https://doi.org/10.1093/qmath/12.1.123; MR0124582; Zbl 0103.05604
[25] J. Sugie, Asymptotic stability of coupled oscillators with time-dependent damping, Qual. Theory Dyn. Syst. 15(2016), No. 2, 553-573. https://doi.org/10.1007/s12346-015-01757; MR3563436; Zbl 1364.34085
[26] G. R. Tomlinson, K. Worden, Nonlinearity in structural dynamics: detection, identification and modelling, Institute of Physics Publishing, Bristol, 2000. Zbl 0990.93508
[27] UnKNOwn Author, https://www.sharetechnote.com/html/DE_Modeling_Example_ SpringMass.html.
[28] Y. XIN, G. Xu, N. Su, Dynamic optimization design of cranes based on human-crane-rail system dynamics and annoyance rate, Shock and Vibration 2017(2017), 8376058. https: //doi.org/10.1155/2017/8376058


[^0]:    ${ }^{\boxtimes}$ Corresponding author.
    Emails: morosanu@math.ubbcluj.ro (G. Moroșanu), cristian.vladimirescu@edu.ucv.ro (C. Vladimirescu).

